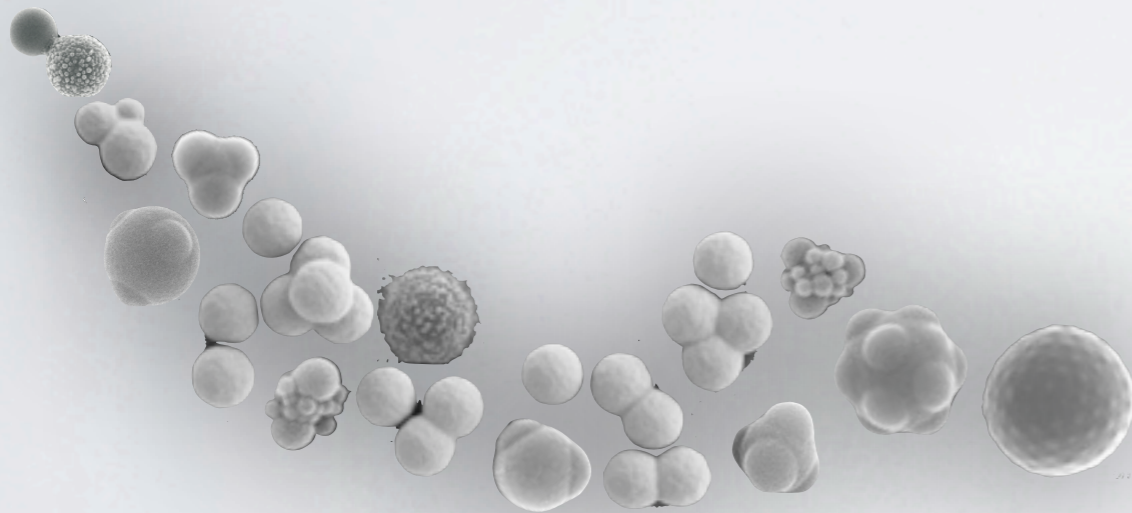


COLLOIDAL SELF-ASSEMBLY



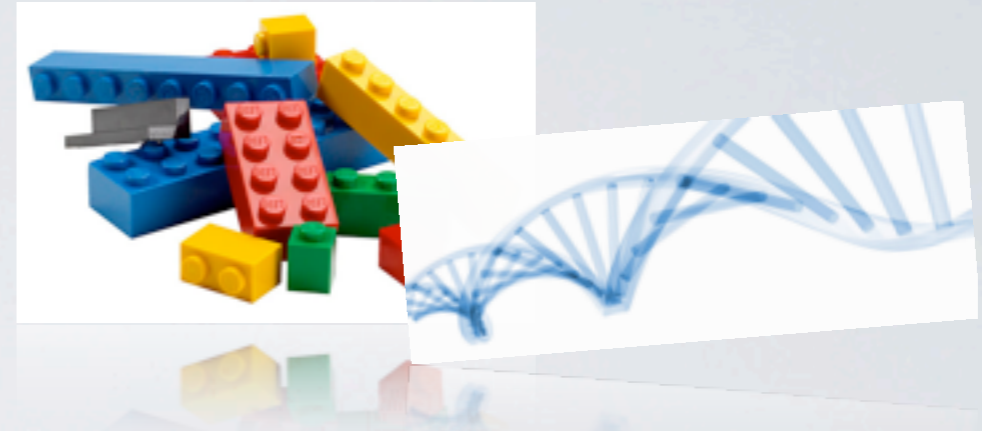
Dr. Daniela J. Kraft



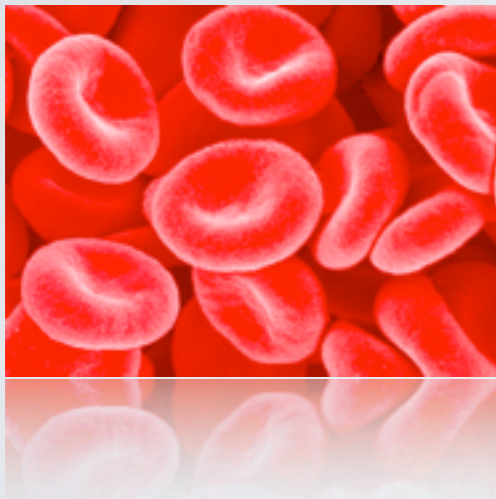
Liquids



Polymers / DNA



Biophysics



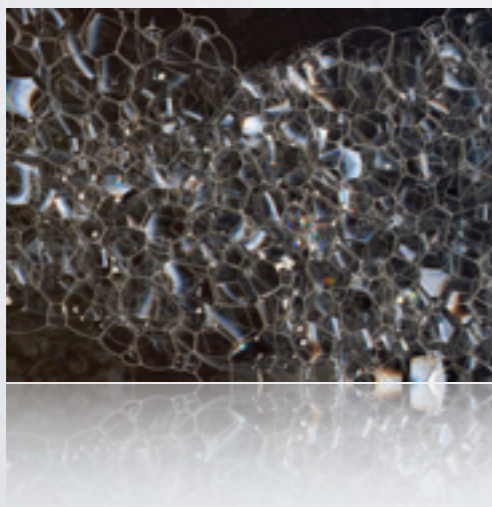
SOFT MATTER

Relevant energy scale $\sim kT$
Easily deformed

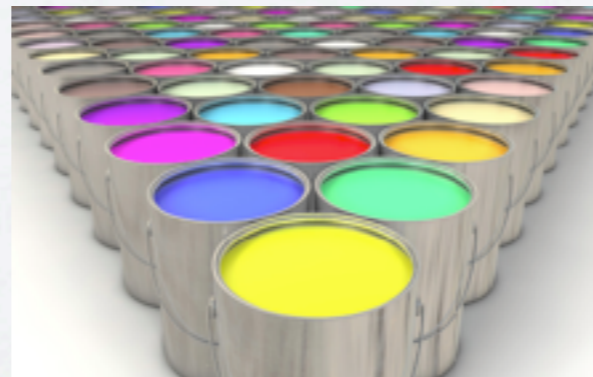
Gels



Foams



Colloids



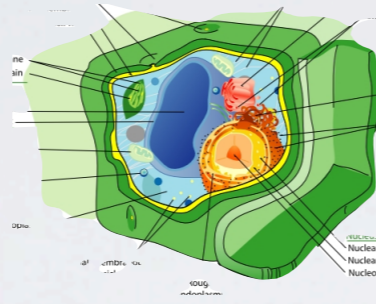
WHAT ARE COLLOIDS?

Planets



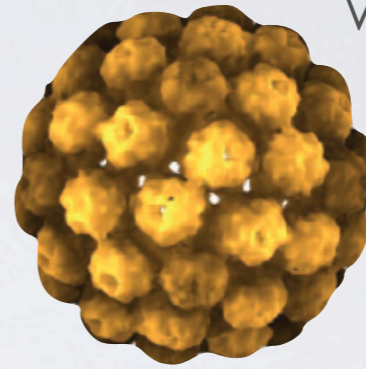
10^5

Cells



10^{-5}

Viruses



10^{-8}

10^{21}

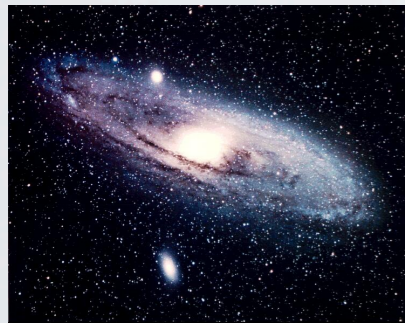
10^0

$10^{-5} \dots 10^{-9}$

10^{-10}

meter

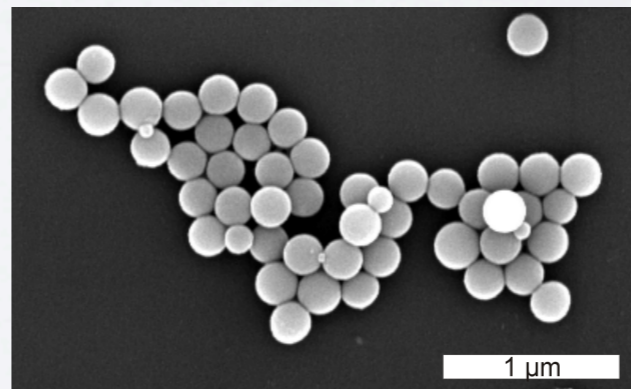
Galaxies



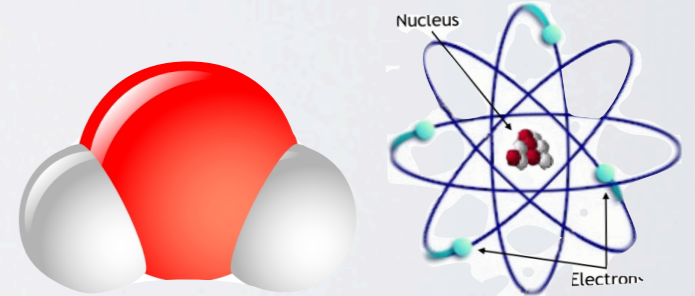
Humans



Colloids



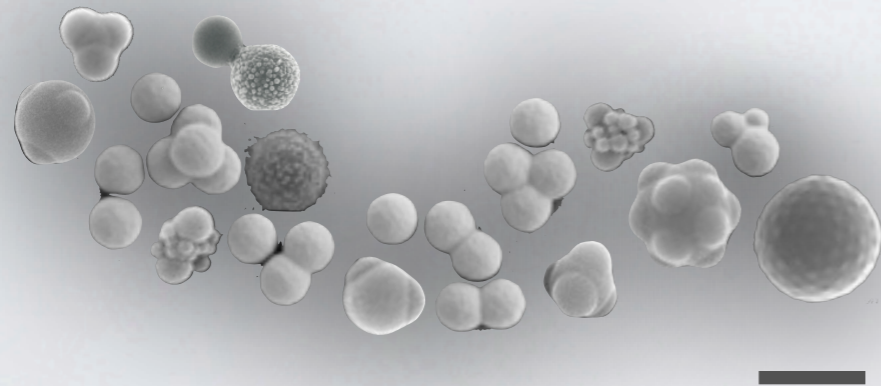
Atoms & molecules



WHAT ARE COLLOIDAL PARTICLES?

Colloids

Small particles suspended in a liquid



Size: $\sim 2\text{nm} - 10\mu\text{m}$

Concentration: $\sim 0.1\%$ to 70% by volume

Materials:

- Plastic: polystyrene, PMMA...
- Inorganic: silica (SiO_2), titania (TiO_2), ...
- Semiconductor: CdSe,...
- Metal: Au, Ag, ...
- Fat, protein,...
- Droplets (emulsions)

Where to find: milk, butter, mayonnaise, blood, ink, paint, toothpaste, coffee,...

WHY ARE WE INTERESTED IN COLLOIDS?

EXPERIMENT: LATEX PAINT STAIN

WHO IS THAT?



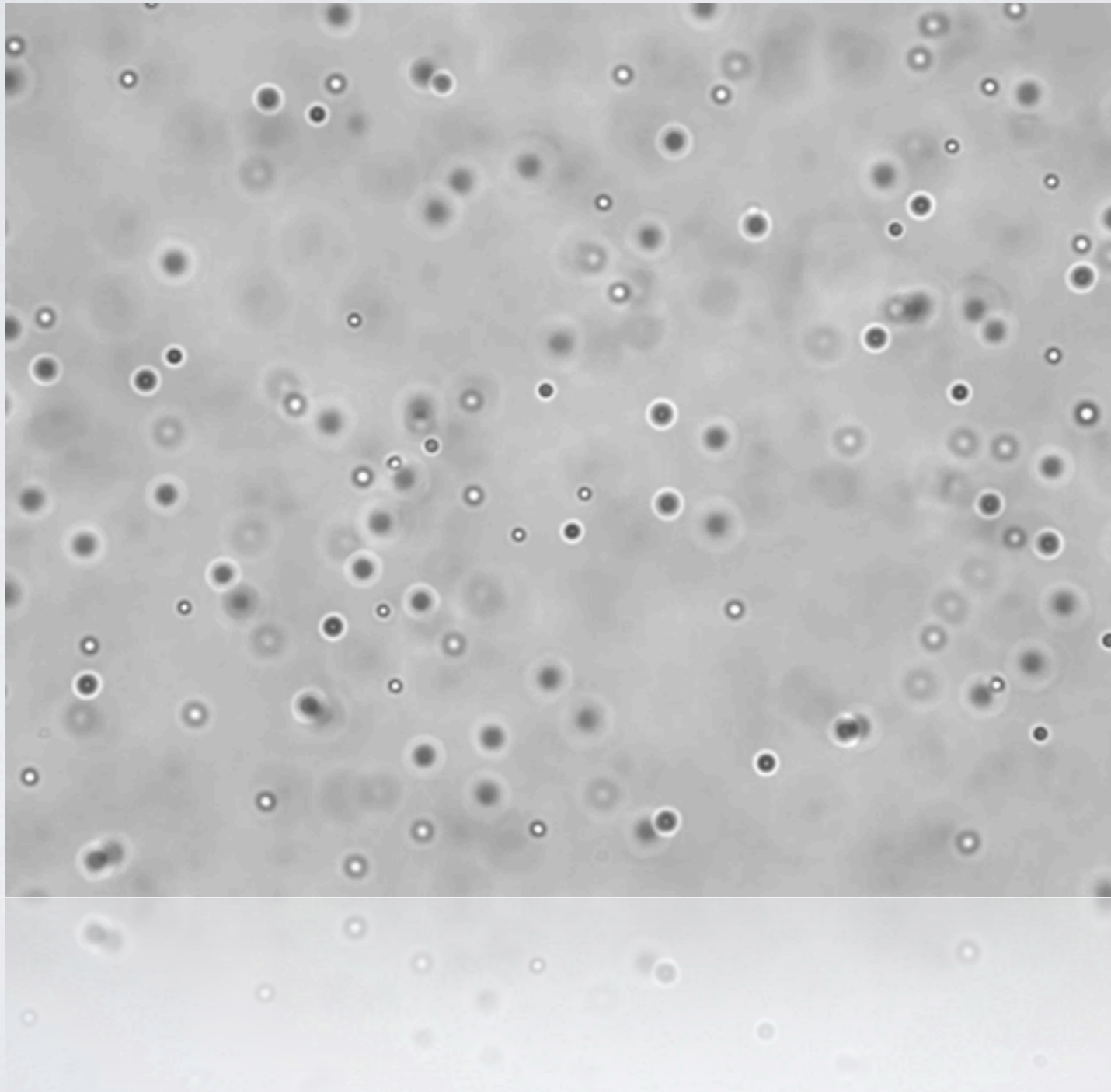
Albert Einstein

Robert Brown

Jean-Baptiste Perrin

Hendrik Casimir

COLLOIDS UNDERGO BROWNIAN MOTION



2 μ m sized particles

WHO IS THAT SCIENTIST?

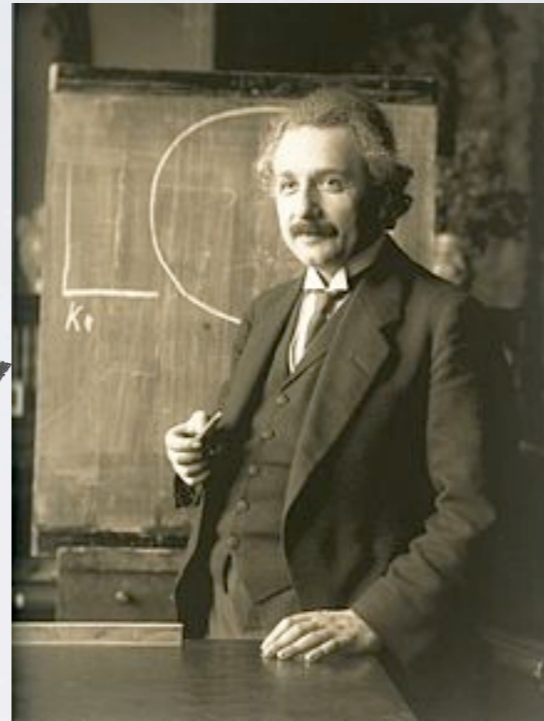


Albert Einstein

Robert Brown

Jean-Baptiste Perrin

Marian Smoluchowski



ANNALEN DER PHYSIK, 17, 549 (1905)

*5. Über die von der molekularkinetischen Theorie
der Wärme geforderte Bewegung von in ruhenden
Flüssigkeiten suspendierten Teilchen;
von A. Einstein.*

Vom Standpunkte der molekularkinetischen Wärmetheorie aus kommt man aber zu einer anderen Auffassung. Nach dieser Theorie unterscheidet sich ein gelöstes Molekül von einem suspendierten Körper *lediglich* durch die Größe, und man sieht nicht ein, warum einer Anzahl suspendierter Körper nicht derselbe osmotische Druck entsprechen sollte, wie der nämlichen Anzahl gelöster Moleküle. Man wird anzunehmen haben, daß

“According to this theory, a dissolved molecule is differentiated from a suspended body only by its size, and it is not apparent why a number of suspended particles should not produce the same osmotic pressure as the same number of molecules.”

Colloids = Atoms

WHO IS THAT SCIENTIST?

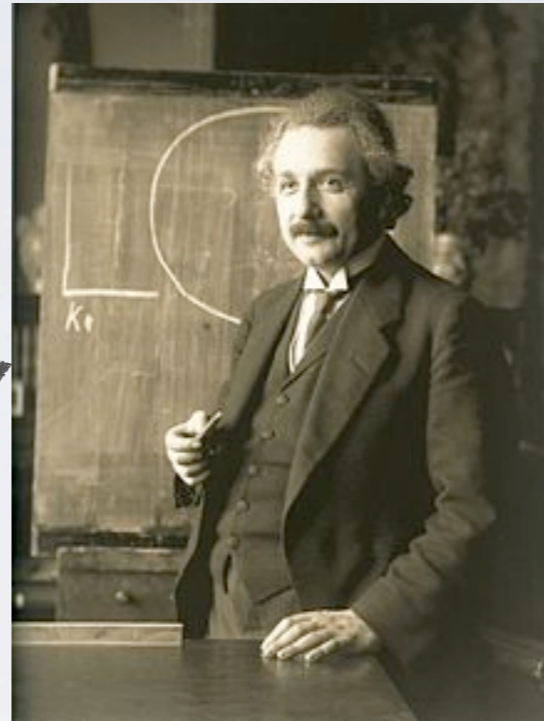


Albert Einstein

Robert Brown

Jean-Baptiste Perrin

Marian Smoluchowski



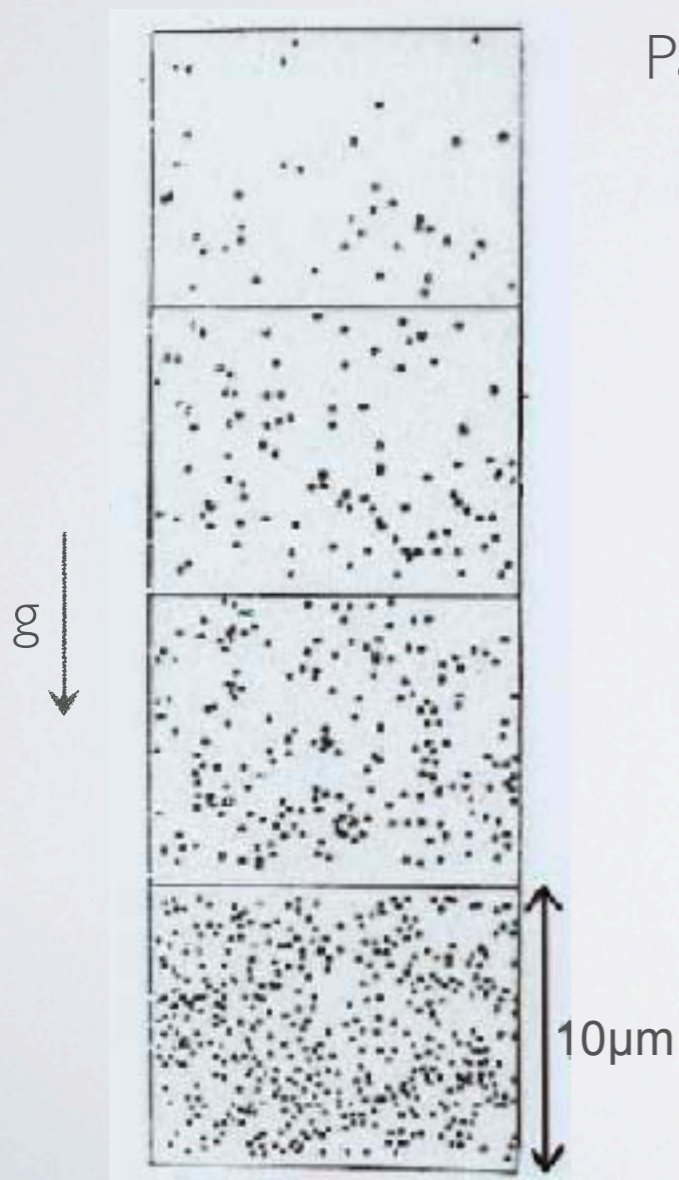
Nobel prize 1926
“For his work on the discontinuous structure of matter, and especially his discovery of the sedimentation equilibrium”

PERRIN'S EXPERIMENT (1907)

Avogadro constant $N_A = 6.02 \cdot 10^{23}$

number of atoms in one gram of hydrogen or 12 g of carbon

Gamboge particles,
R=0.3um



Particles are Boltzmann distributed in the gravitational potential

$$\frac{n(z)}{n(0)} = \exp \left[- \frac{V(z) - V(0)}{k_B T} \right]$$

$$V(z) = \Delta m g z$$

Archimedes

$$\Delta m = \frac{4\pi}{3} R^3 (\rho_p - \rho_o)$$

$$\frac{n(z)}{n(0)} = \exp \left[- \frac{\Delta m g z}{k_B T} \right]$$

Number of Avogadro

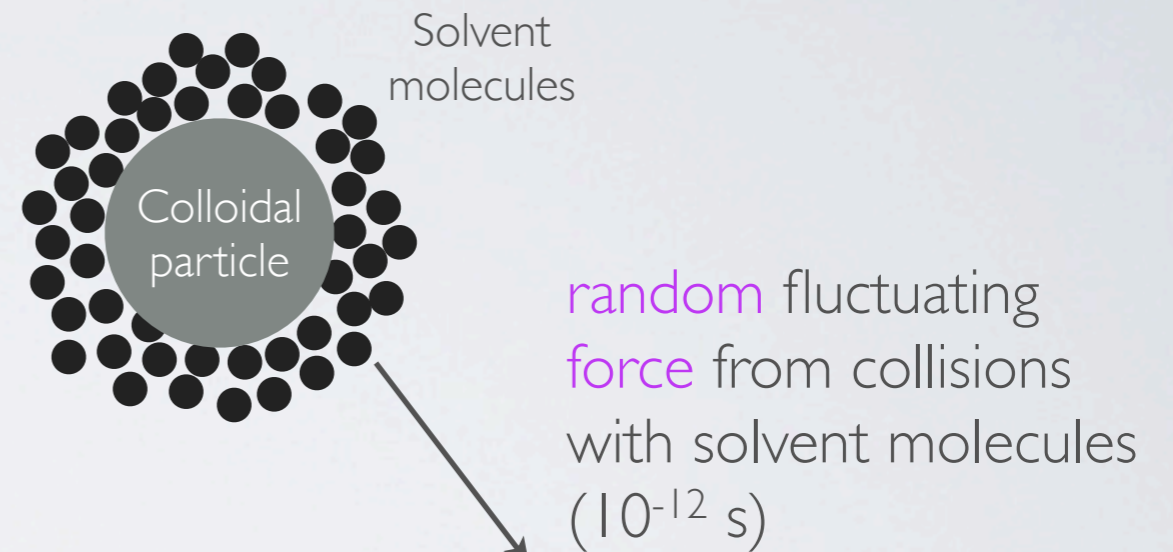
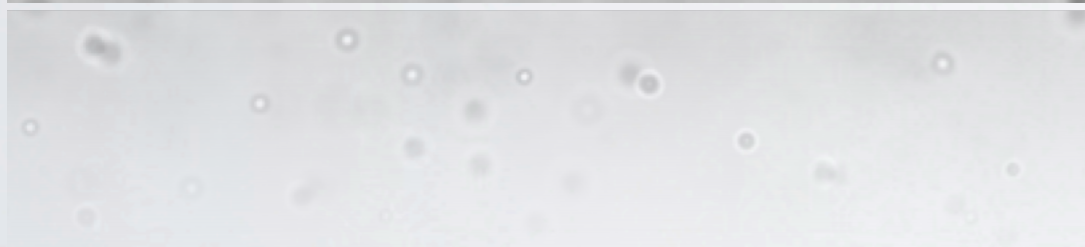
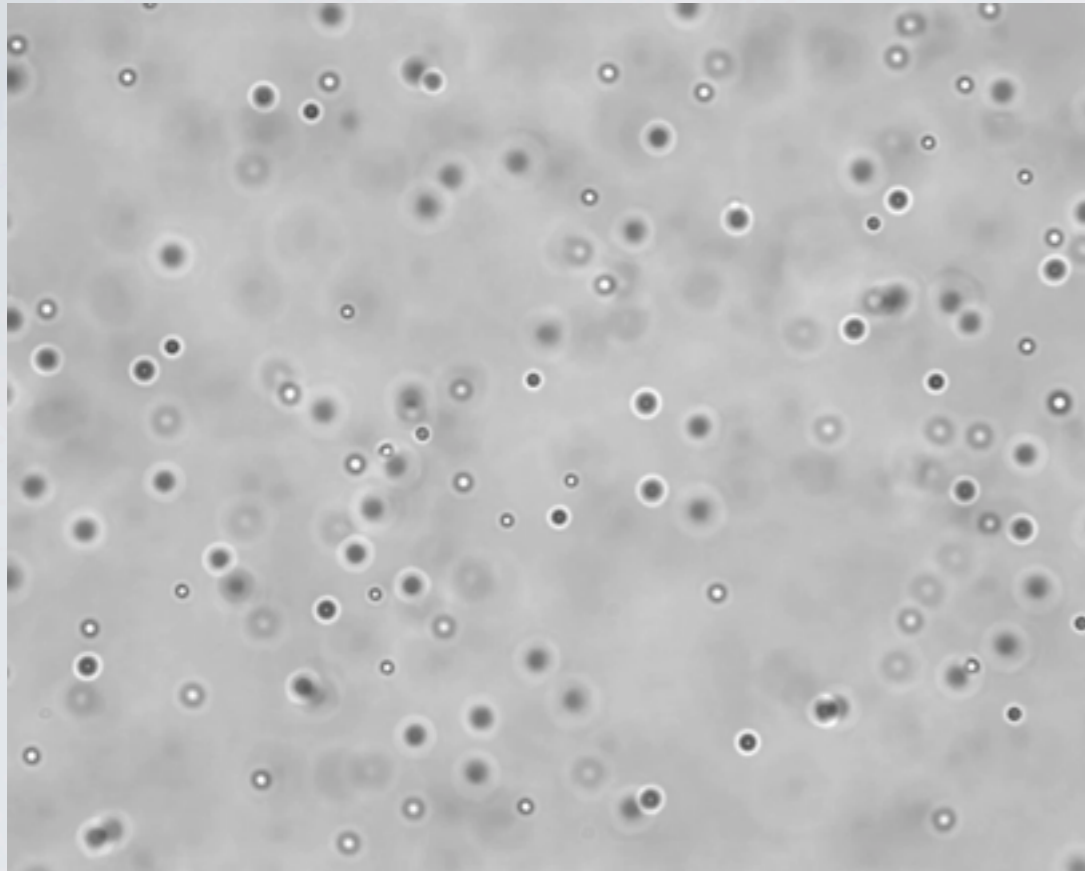
$$N_A = \frac{R}{k_B}$$

Perrin

$$N_A = 6.82 \cdot 10^{23}$$

Colloids = Atoms

COLLOIDS UNDERGO BROWNIAN MOTION



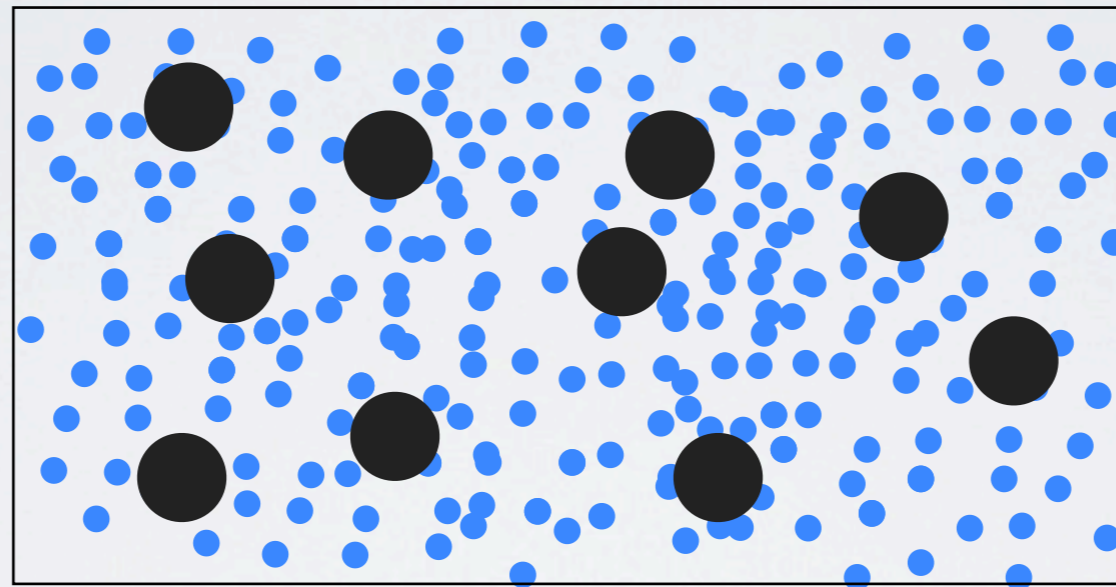
$$m \frac{d\vec{v}(t)}{dt} = -\zeta \vec{v}(t) + \vec{\xi}(t)$$

mass of particle velocity of particle viscous friction

For spheres: $\zeta = 6\pi\eta r$

Brownian motion of colloidal particles is just a large-scale manifestation of the thermal motion of the solvent molecules.

CONCENTRATED COLLOIDAL SUSPENSIONS

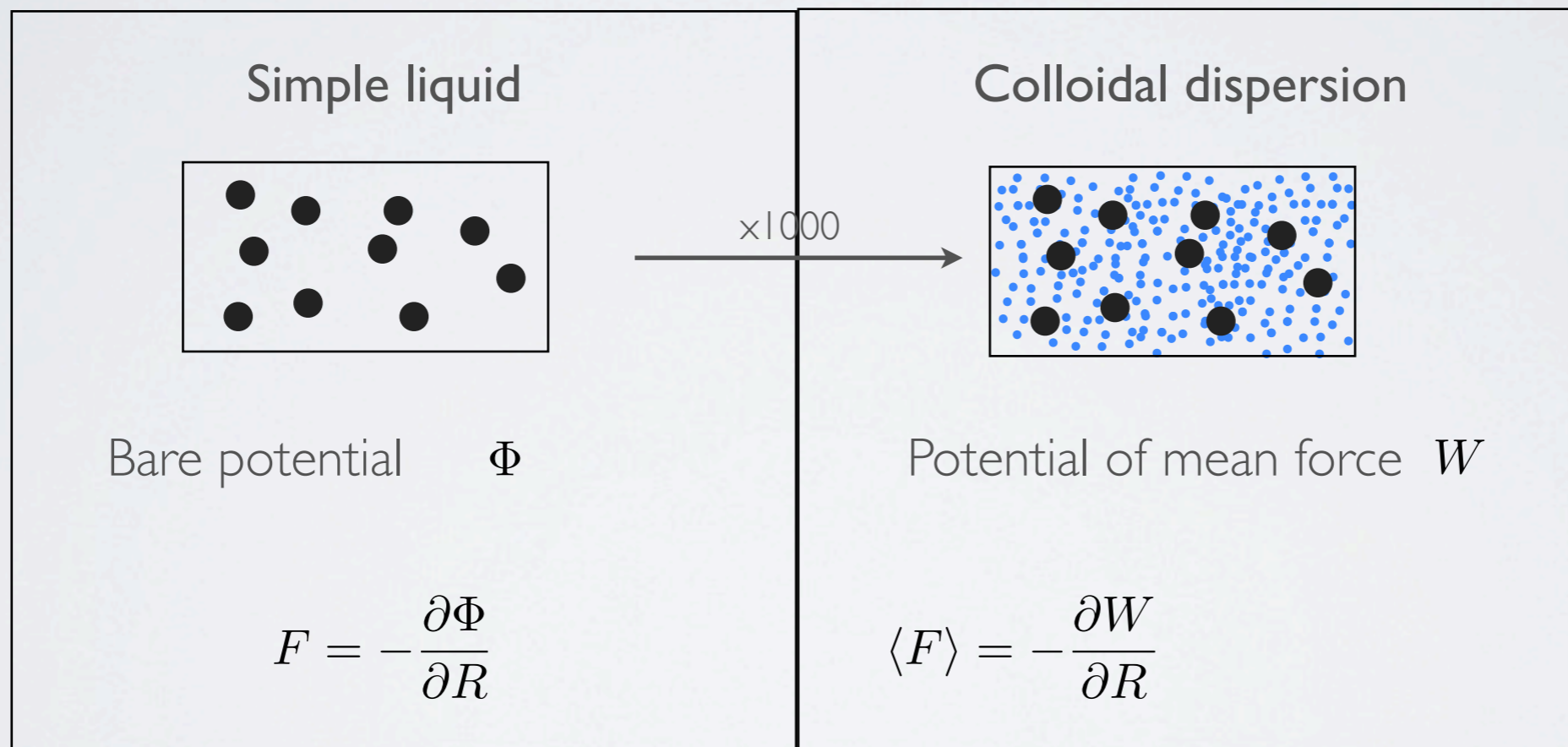


Particle interactions
+
Statistical physics

COLLOID-ATOM ANALOGY

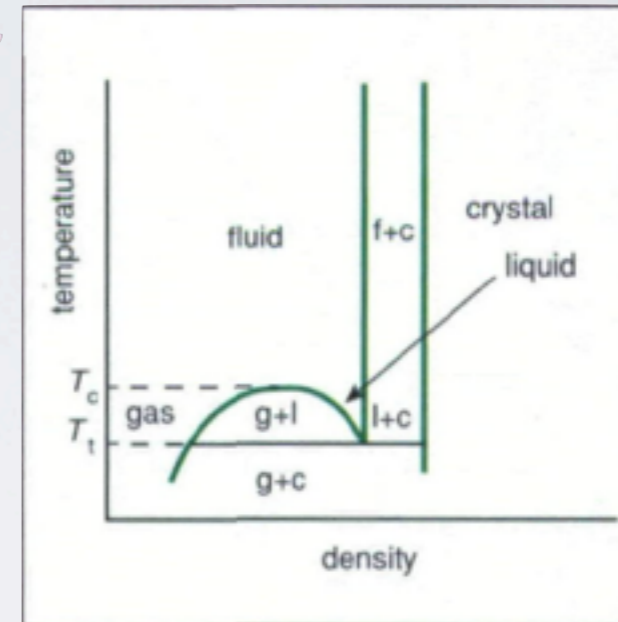
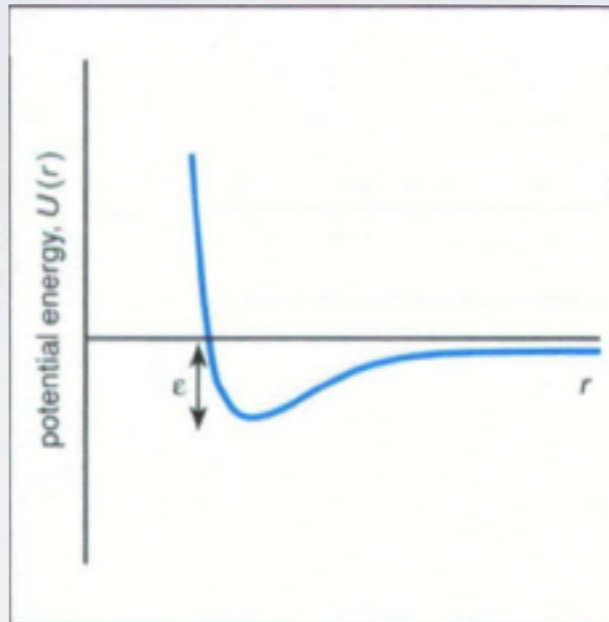
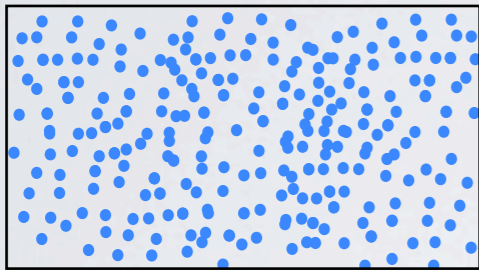
Key concept

The same concepts and methods of the statistical mechanics of simple liquids can be applied to colloidal suspensions. One must simply replace the bare potential $V(r)$ by the potential of the average force $W(r)$.

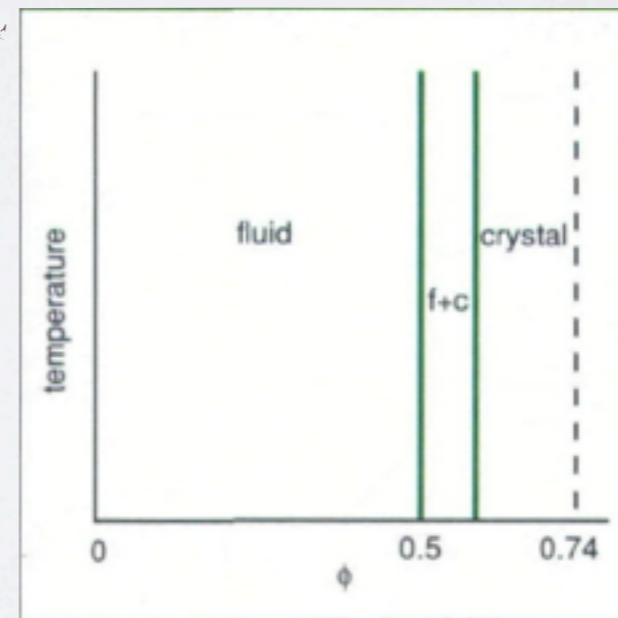
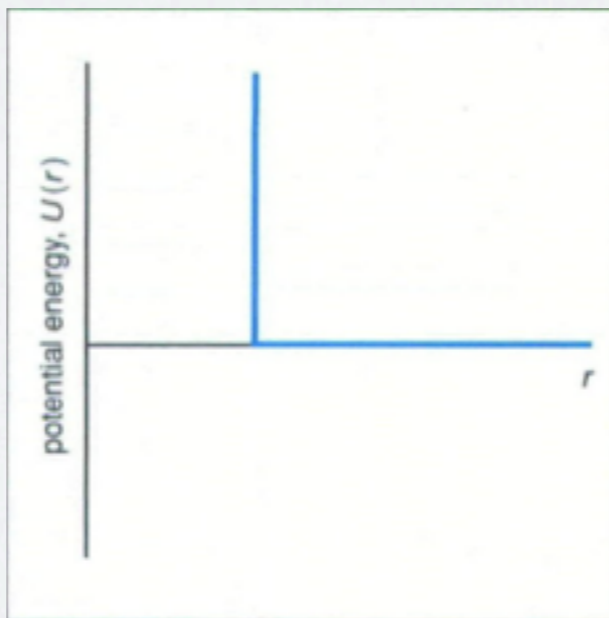
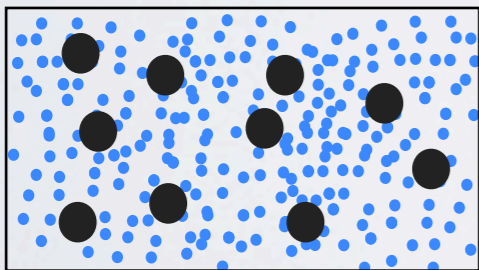


PHASE BEHAVIOR OF ATOMS AND COLLOIDS

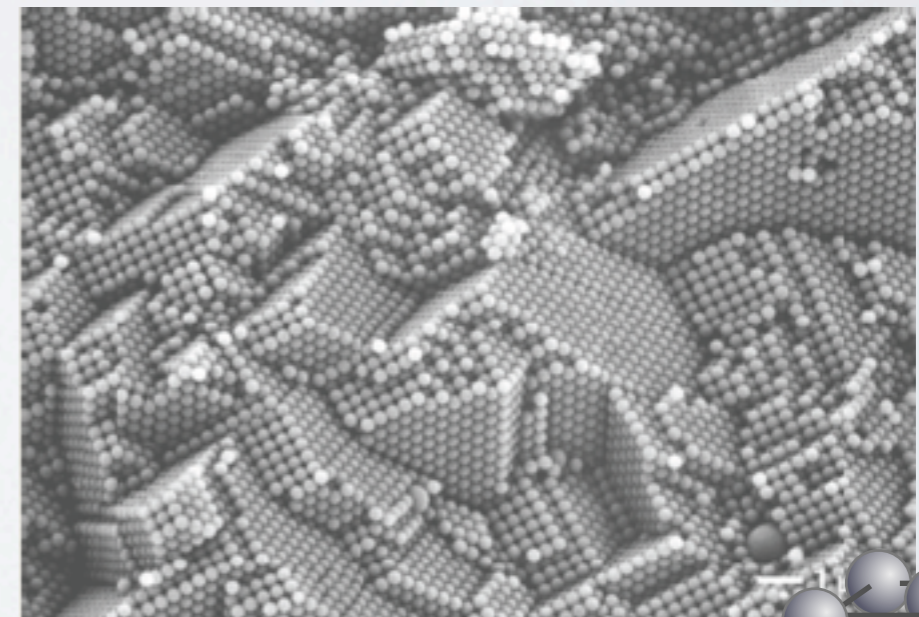
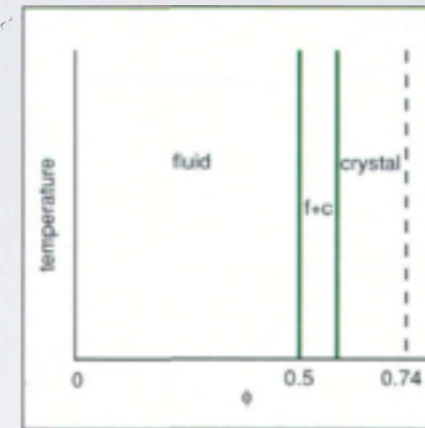
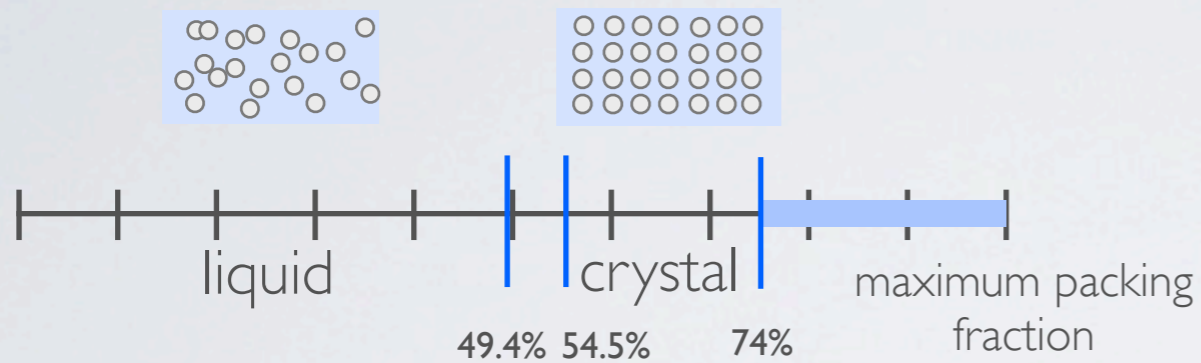
Atoms



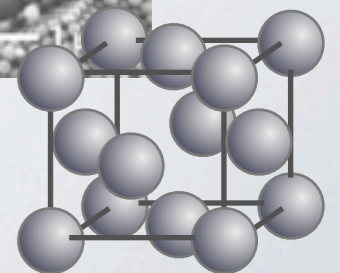
Hard-sphere colloid



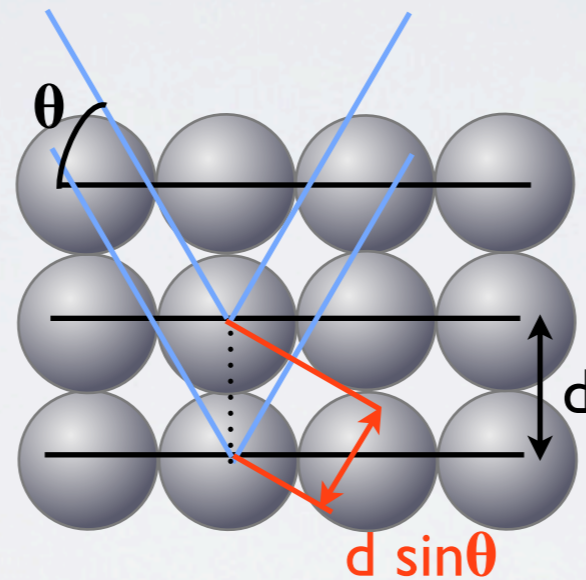
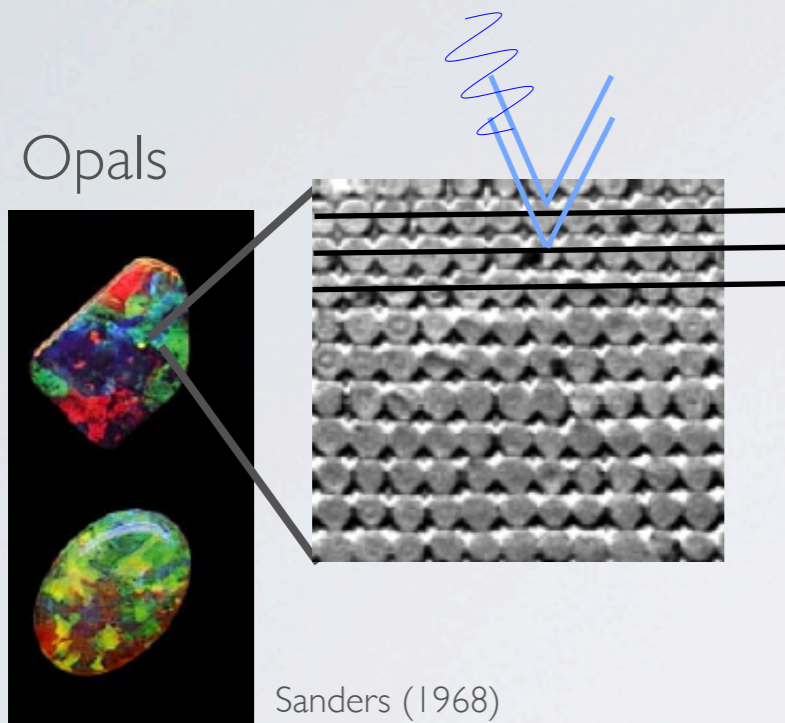
PHASE DIAGRAM OF HARD SPHERE COLLOIDS



Colloidal fcc crystal made with lattice constant 0.5 μ m



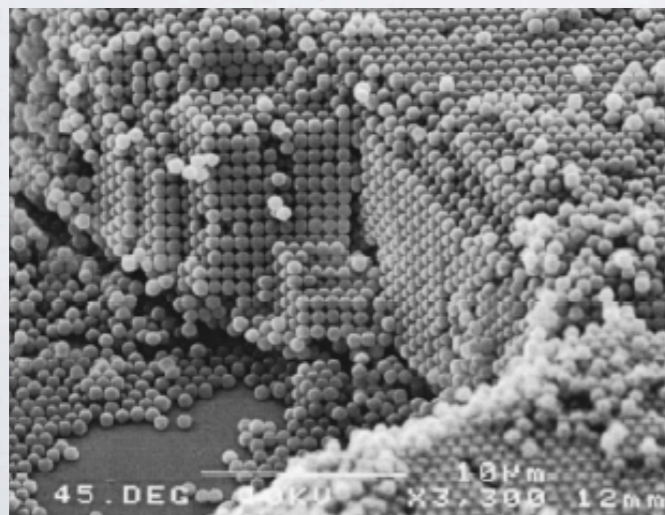
EXPERIMENT: LATEX PAINT STAIN



Bragg's law

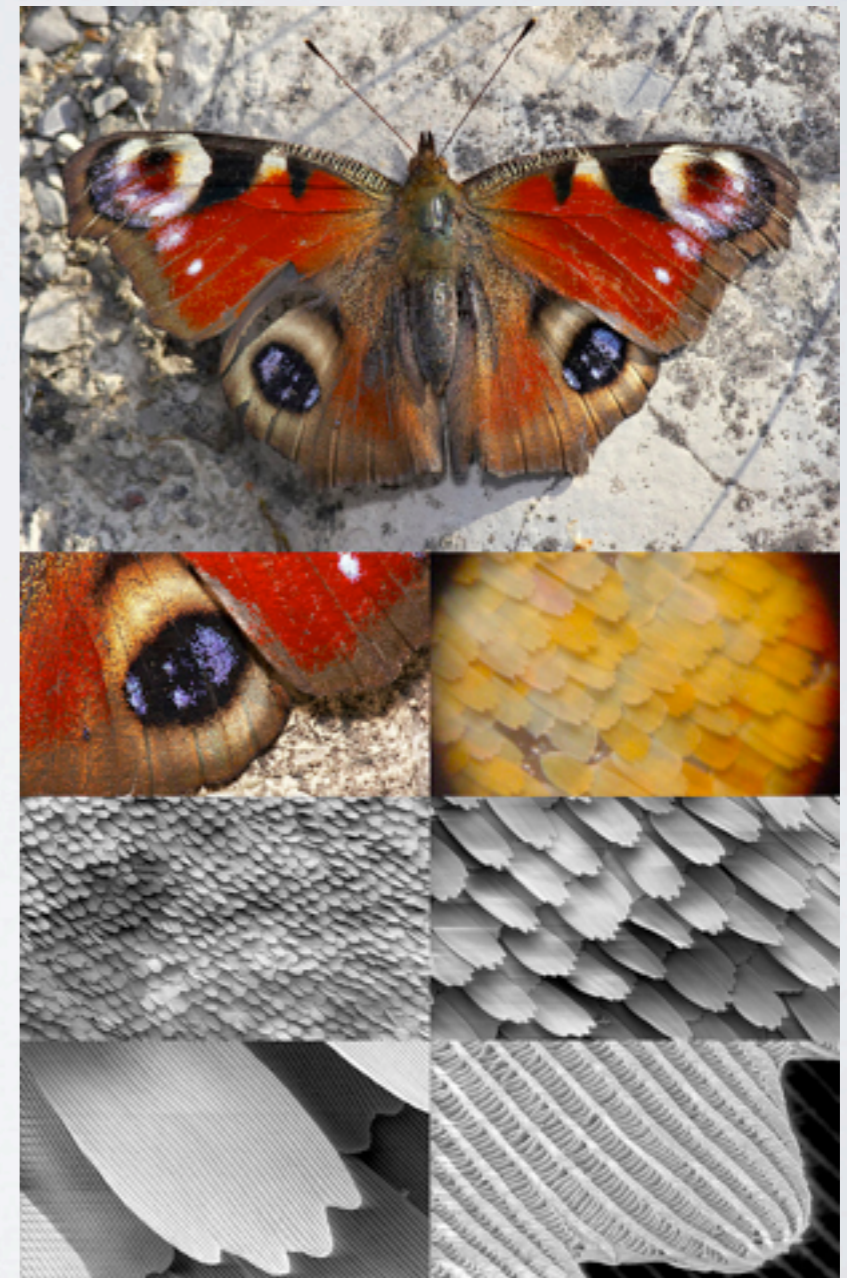
$$2d \sin \theta = n\lambda$$

Colloidal crystal



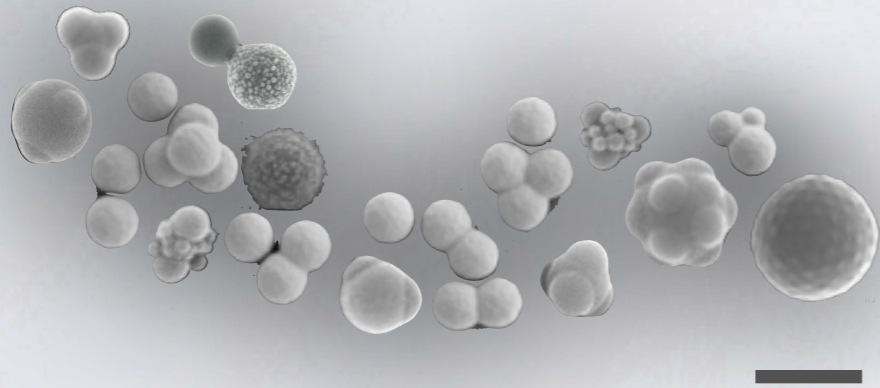
Amos et al., PRE 61, 2929 (2000)

Iridescent feathers, bugs,...



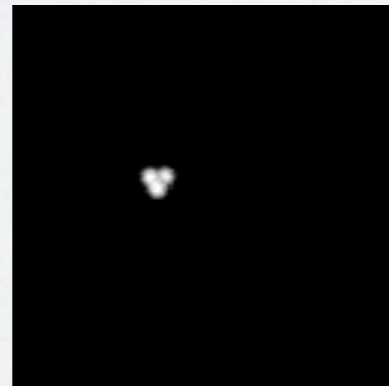
"Butterfly magnification series collage", Wikipedia

WHY COLLOIDAL PARTICLES?



Size:
~2nm-10 μ m

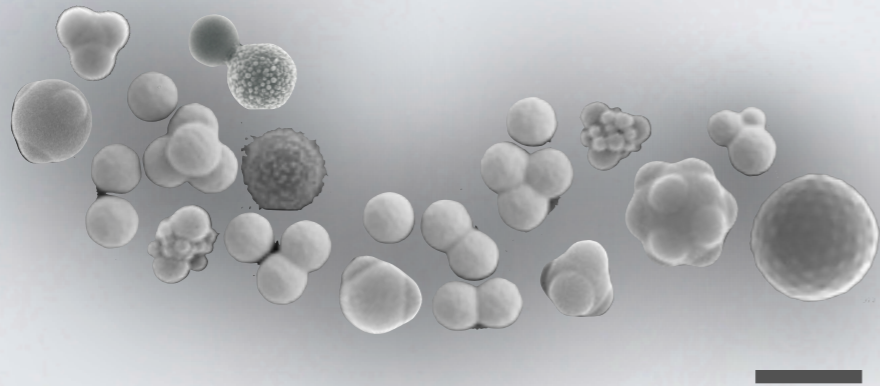
- Constant randomization of the sample
 - System reaches energy minimum
 - Self-assembly
- Observable by microscopy



Kraft et al. PRE 88(2013)

- Tunable interactions

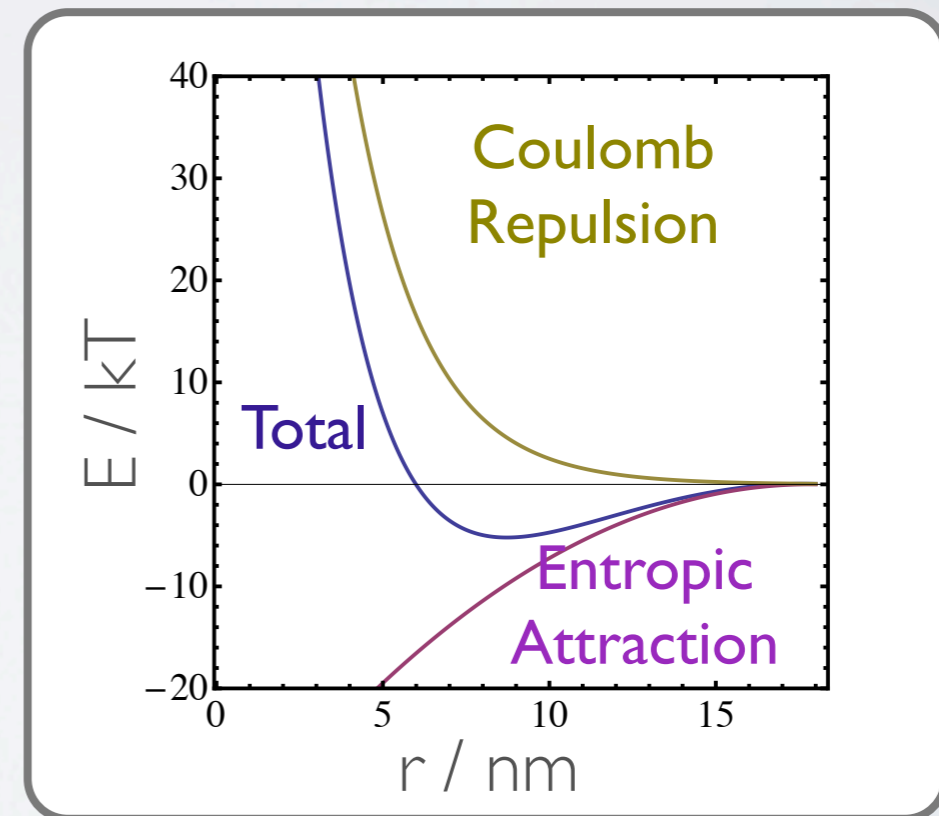
WHY COLLOIDAL PARTICLES?



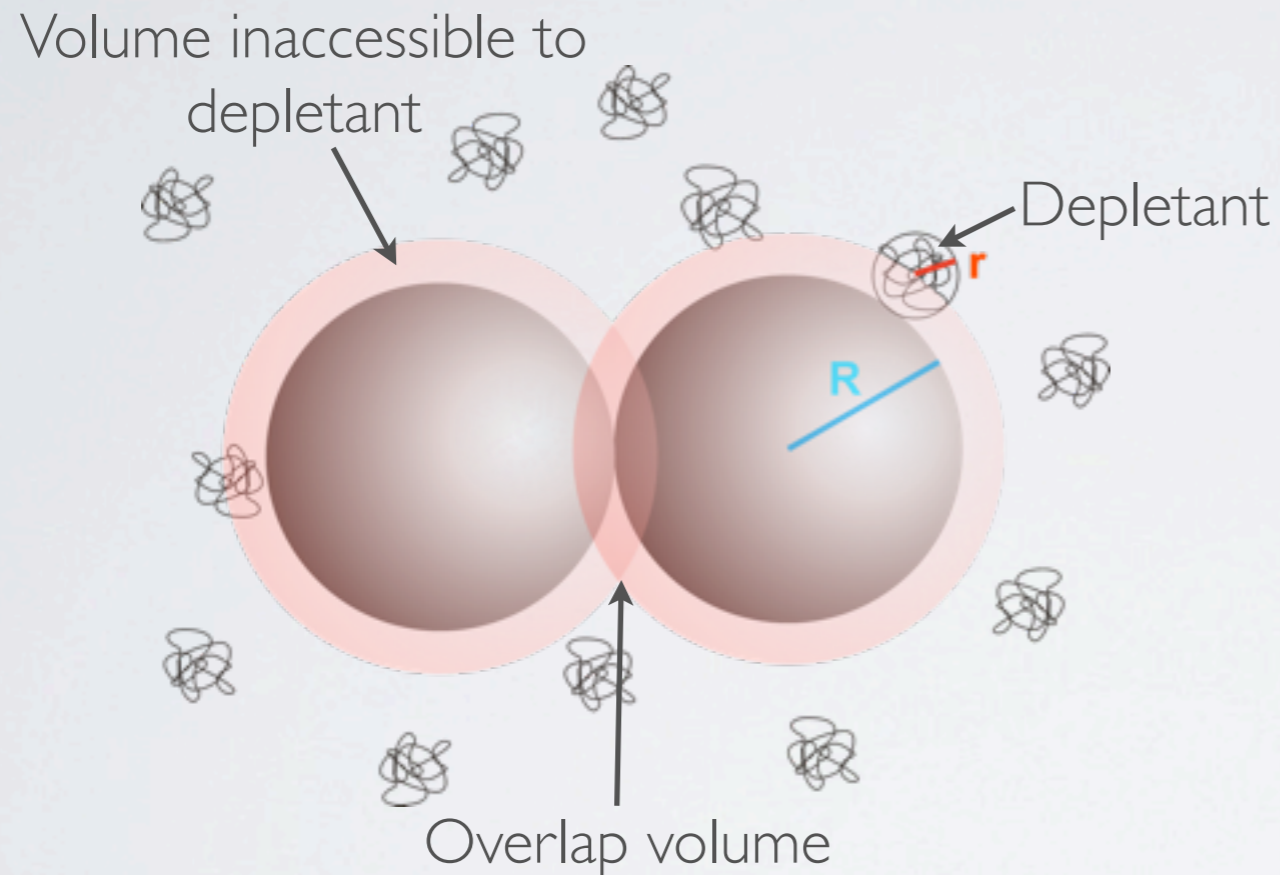
Size:
~2nm-10 μ m

→ Tunable interactions

- Van der Waals forces
- Coulomb Repulsion
- Magnetic forces
- Entropic forces (depletion)
- Sticky DNA linkers
- Steric repulsion
-



THE DEPLETION INTERACTION



Effective attraction between larger colloid!

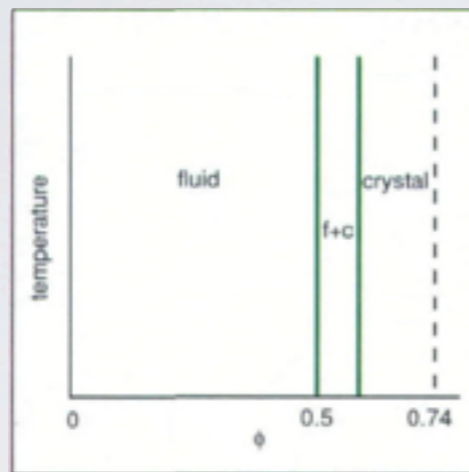
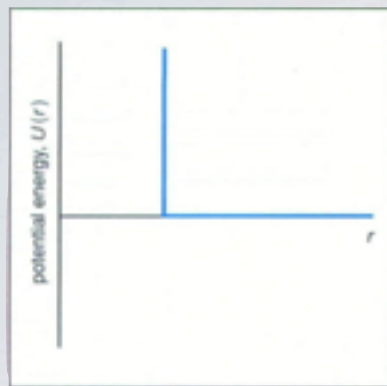
$$\begin{aligned} u &= -\Pi V_{overlap} \\ &= -k_B T n_{depl} V_{overlap} \end{aligned}$$

Depletion forces are entropic forces!

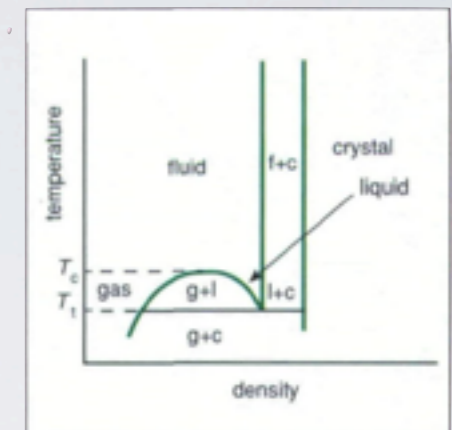
TUNABLE INTERACTIONS

TUNABLE PHASE DIAGRAM

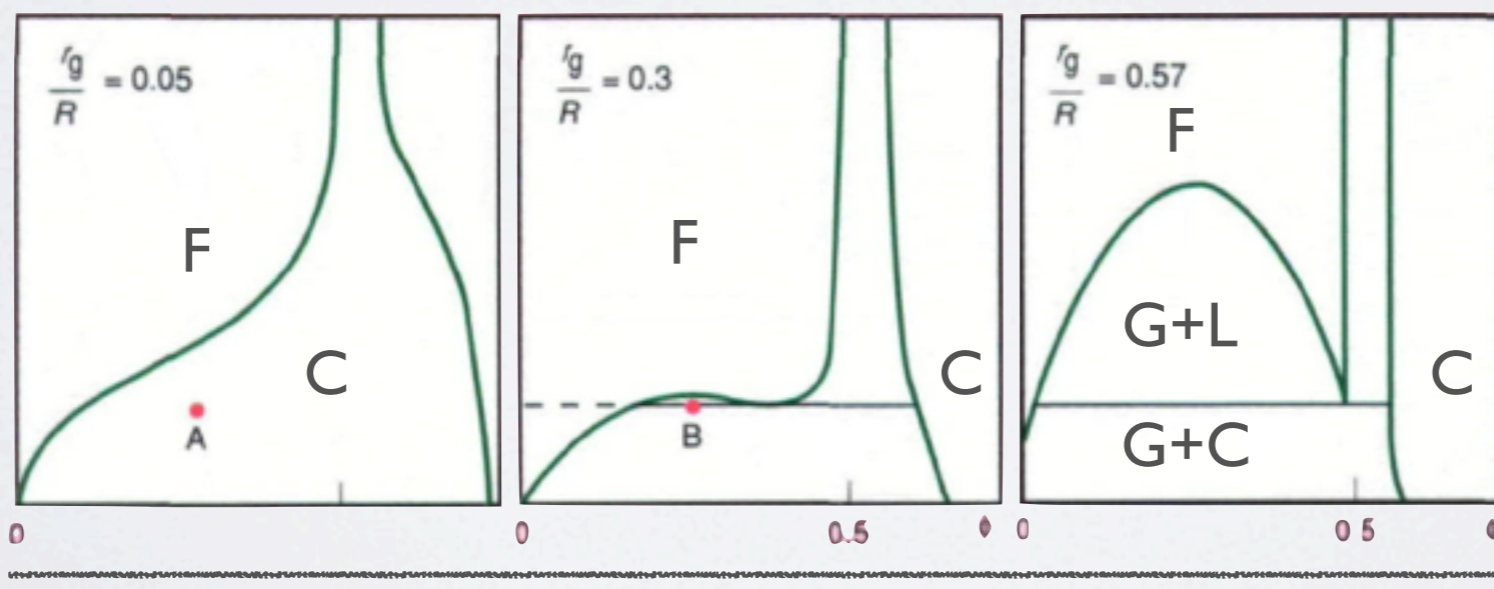
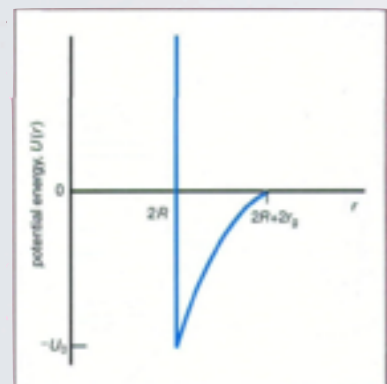
Hard sphere interaction



atom phase diagram

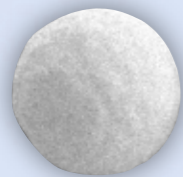


Hard sphere + depletion attraction

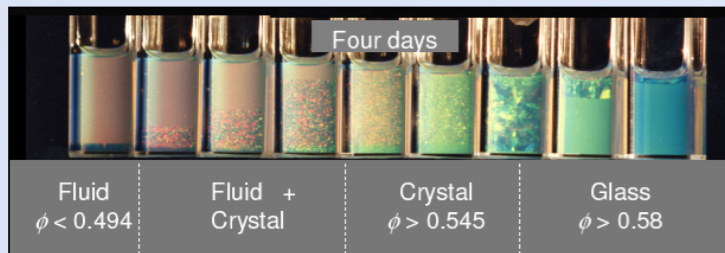


polymer size / colloid size
= interaction range

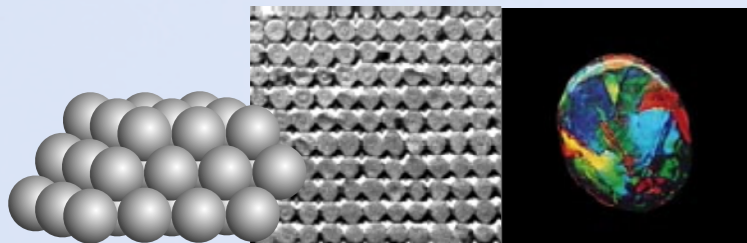
FROM SPHERES TO COMPLEX PARTICLES



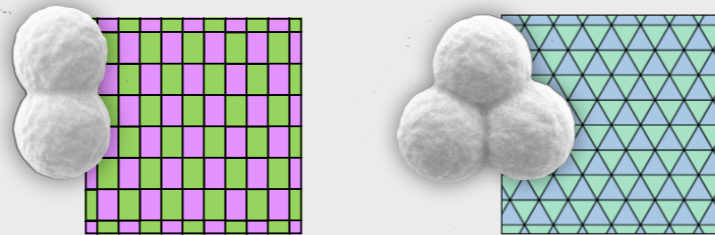
model 'atom'



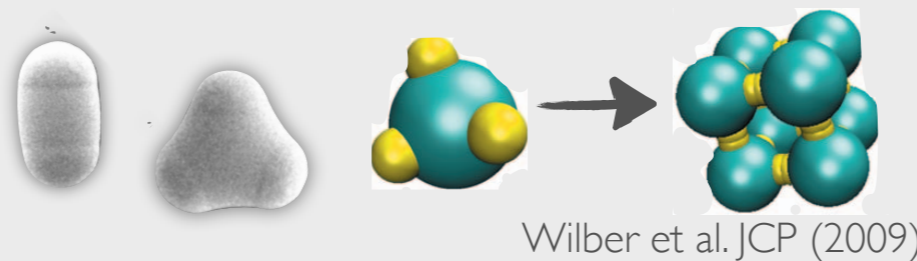
Pusey, van Meegen, Nature (1986)



Anisotropic shape



Anisotropic interactions



Highly specific interactions



+ External guiding rules
+ activity

Particle shape and interactions

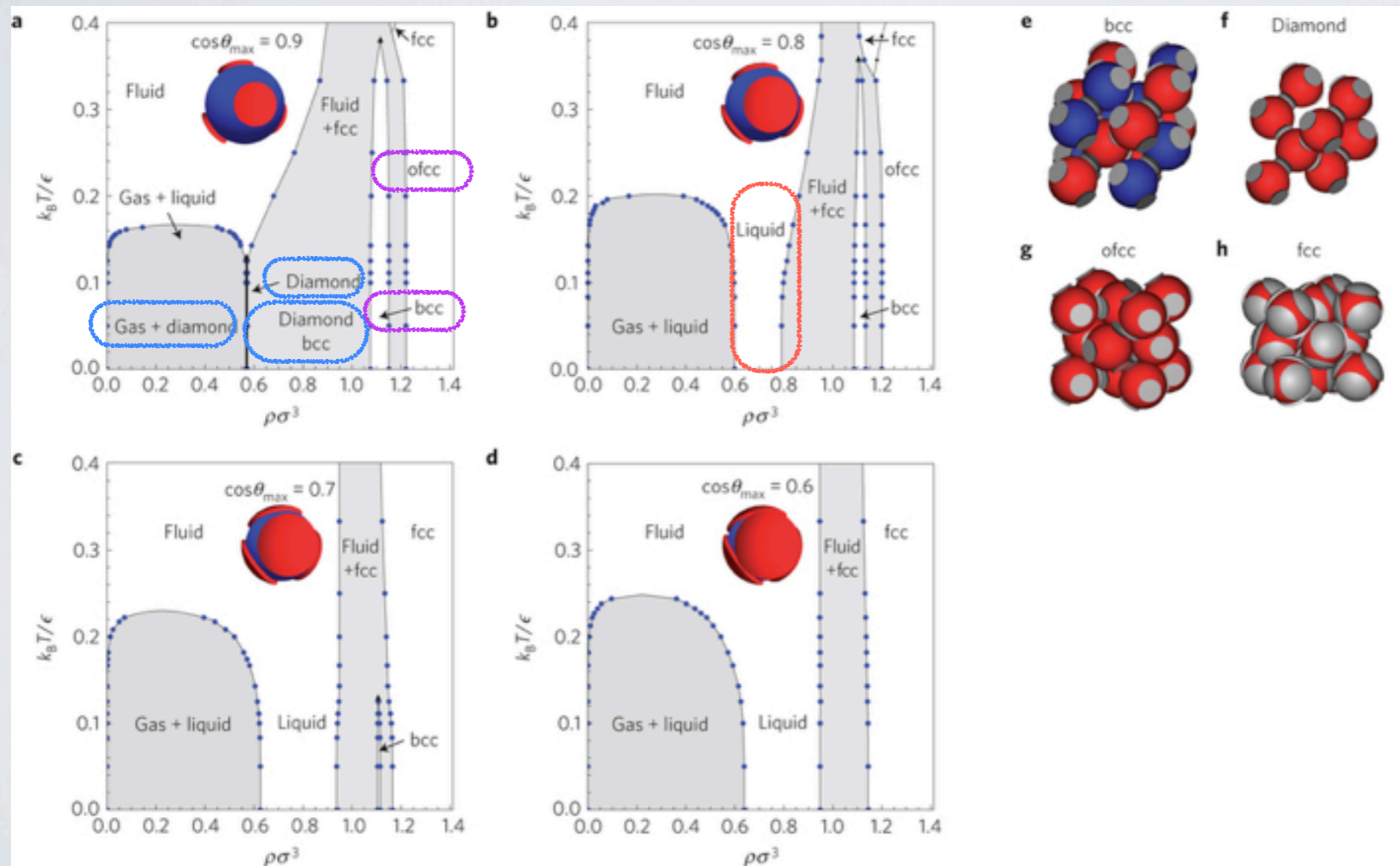


Assembled structure

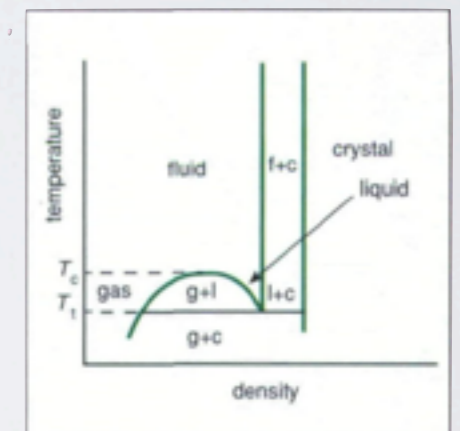


Understand & Design self-assembly

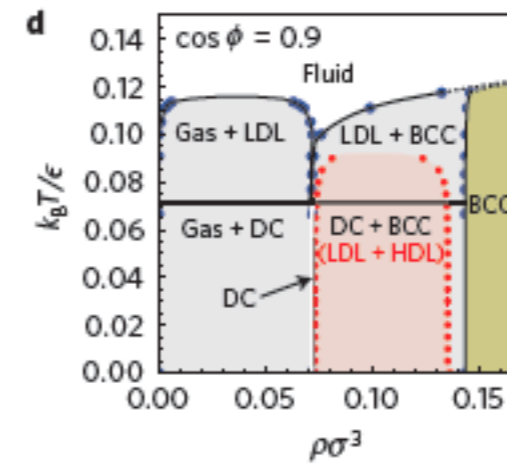
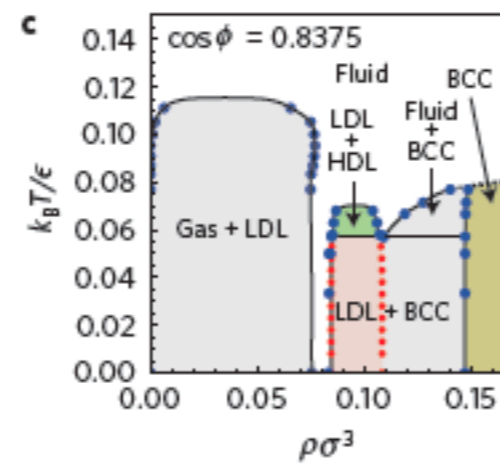
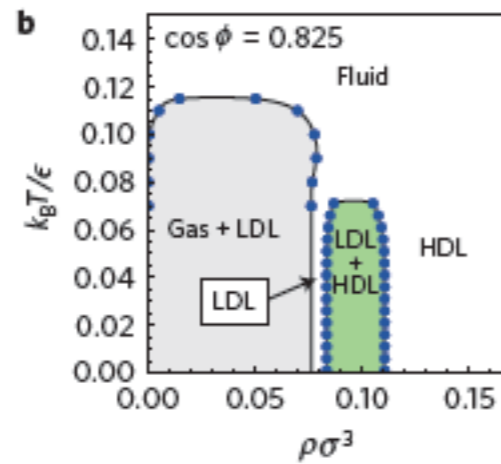
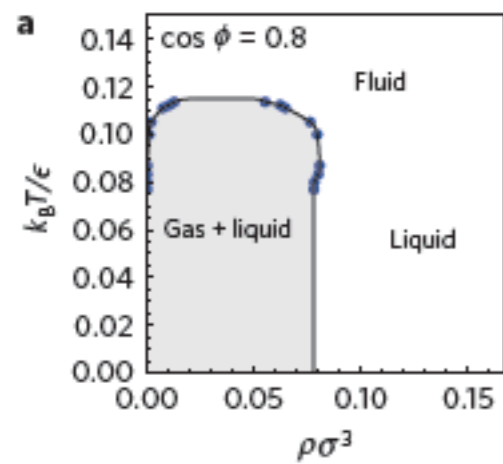
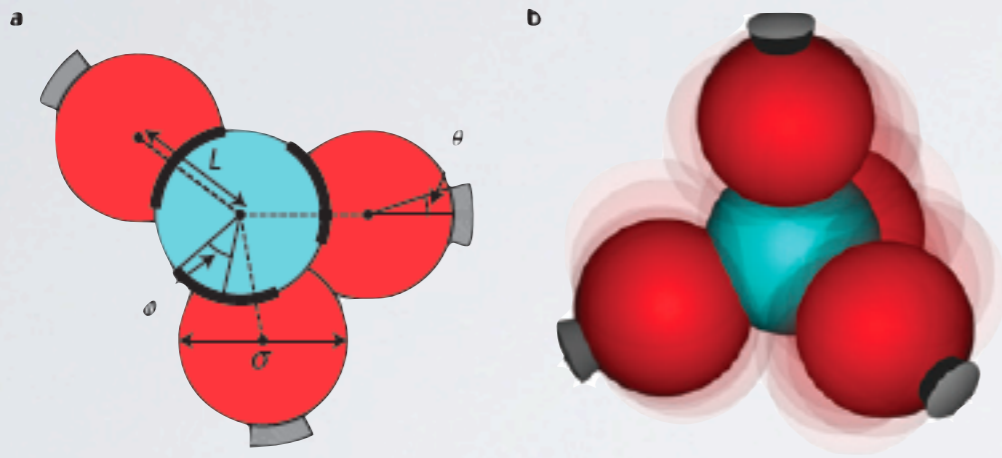
PATCHY PARTICLES: LIQUID PHASE CAN BE PREFERRED OVER THE SOLID PHASE!



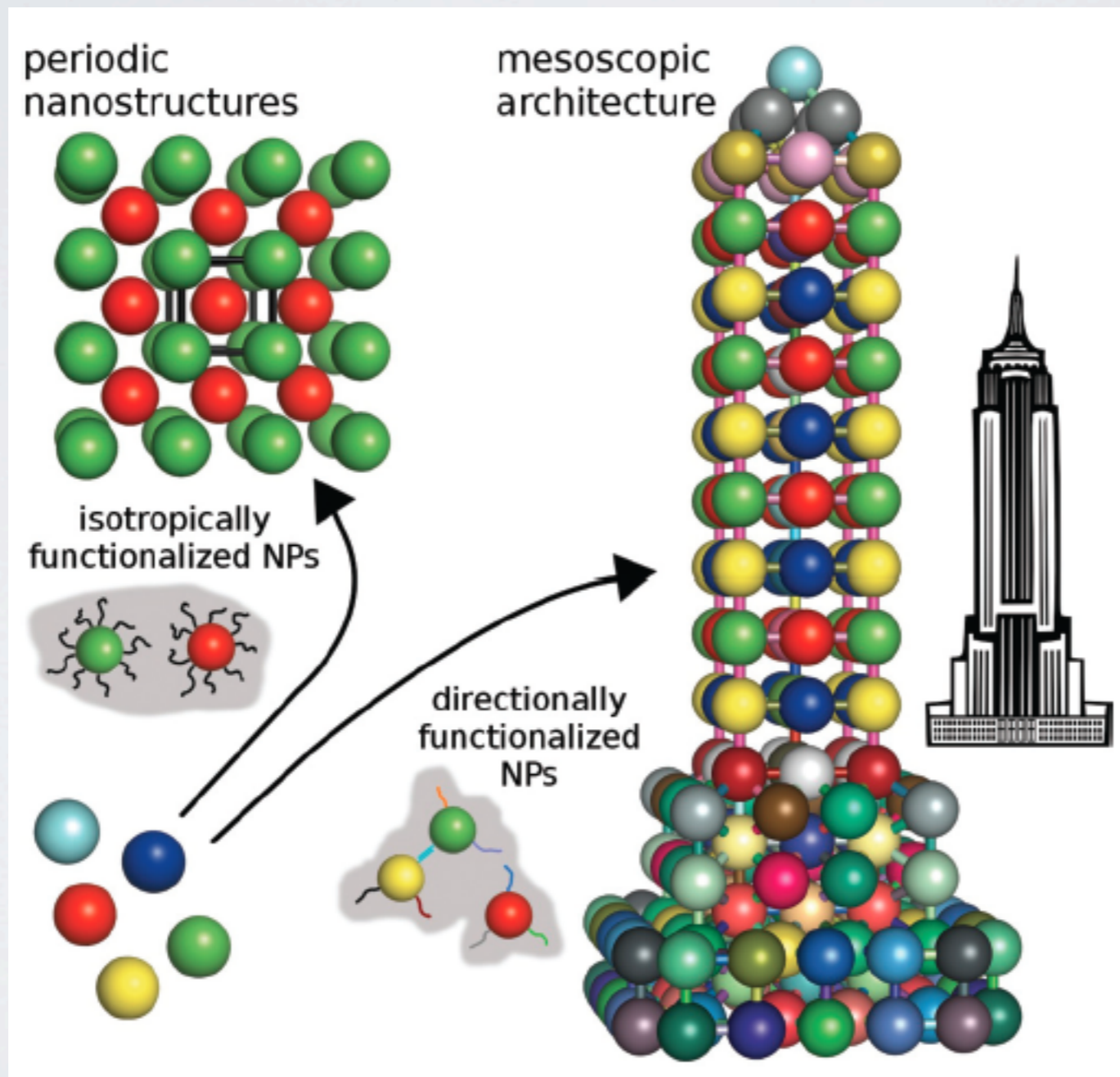
simple liquid phase diagram



FLEXIBLE AND DIRECTIONAL INTERACTIONS MAY LEAD TO STRANGE NEW PHASE BEHAVIOR



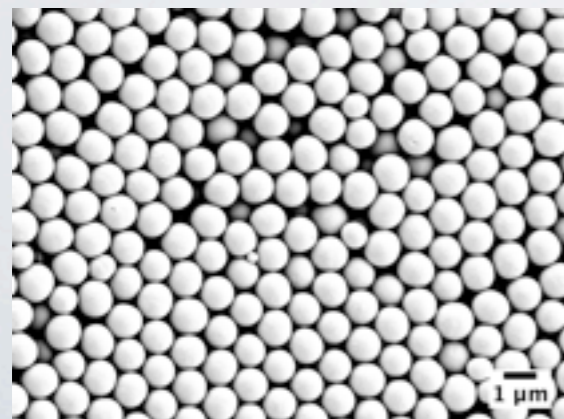
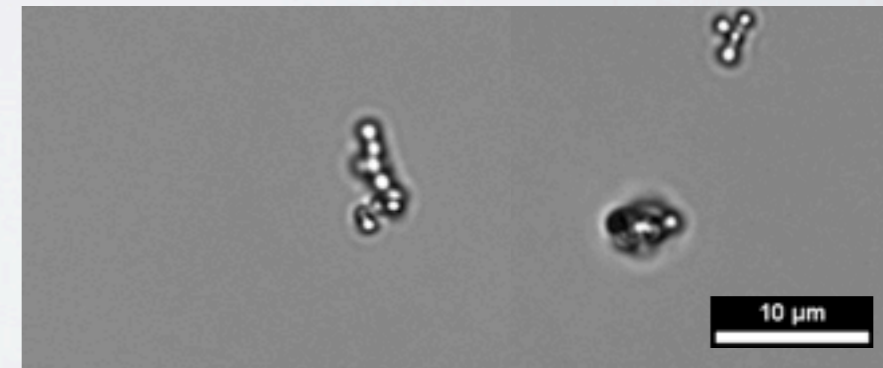
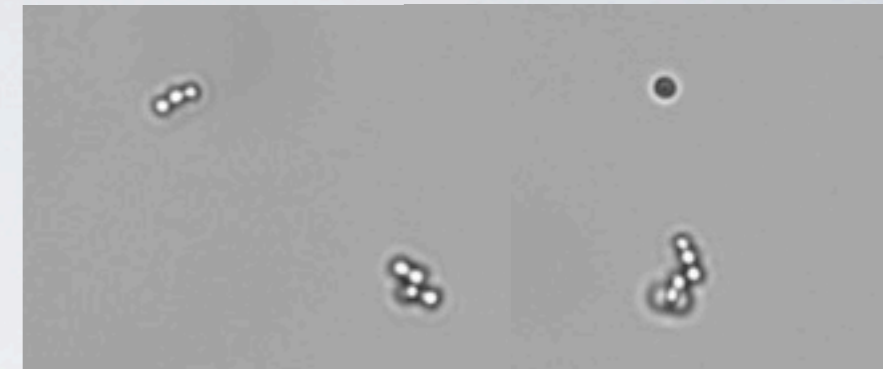
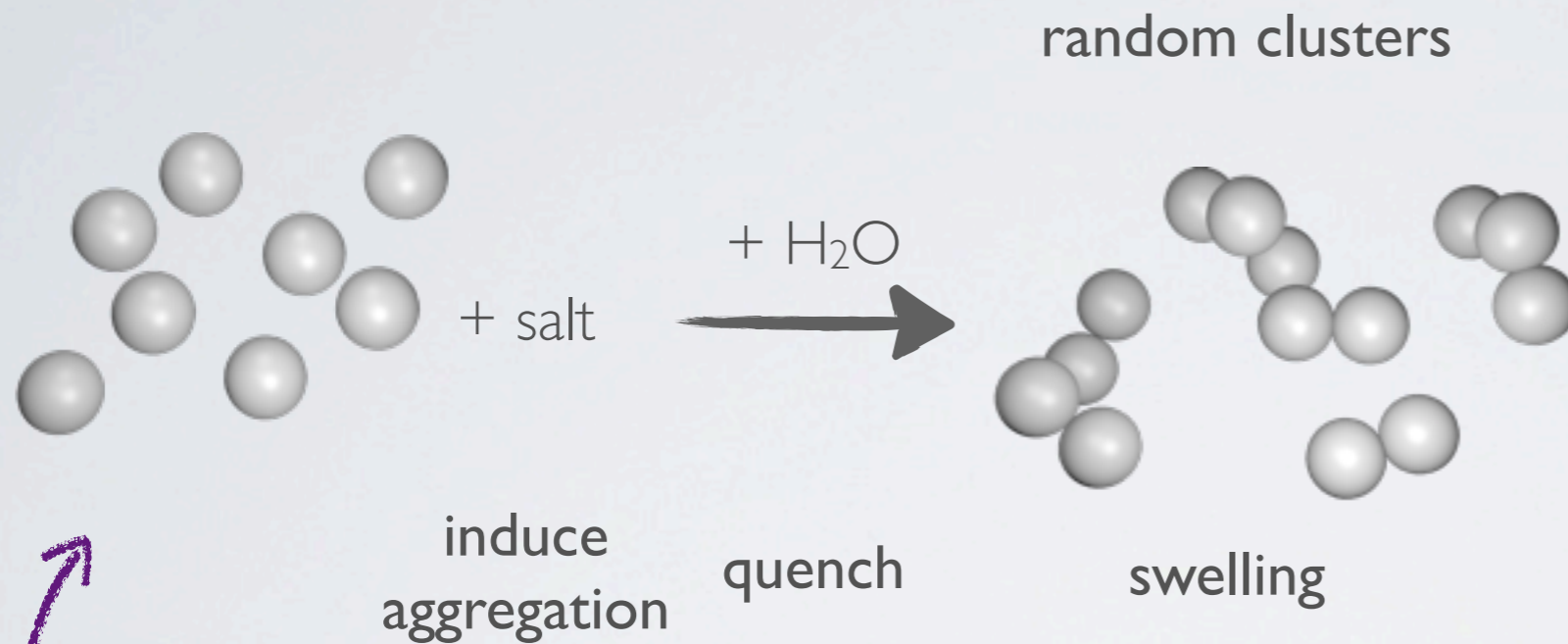
DNA-COATED COLLOIDS CAN IN PRINCIPLE BUILD ARBITRARILY COMPLEX STRUCTURES





HOW TO MAKE NON-SPHERICAL PARTICLES

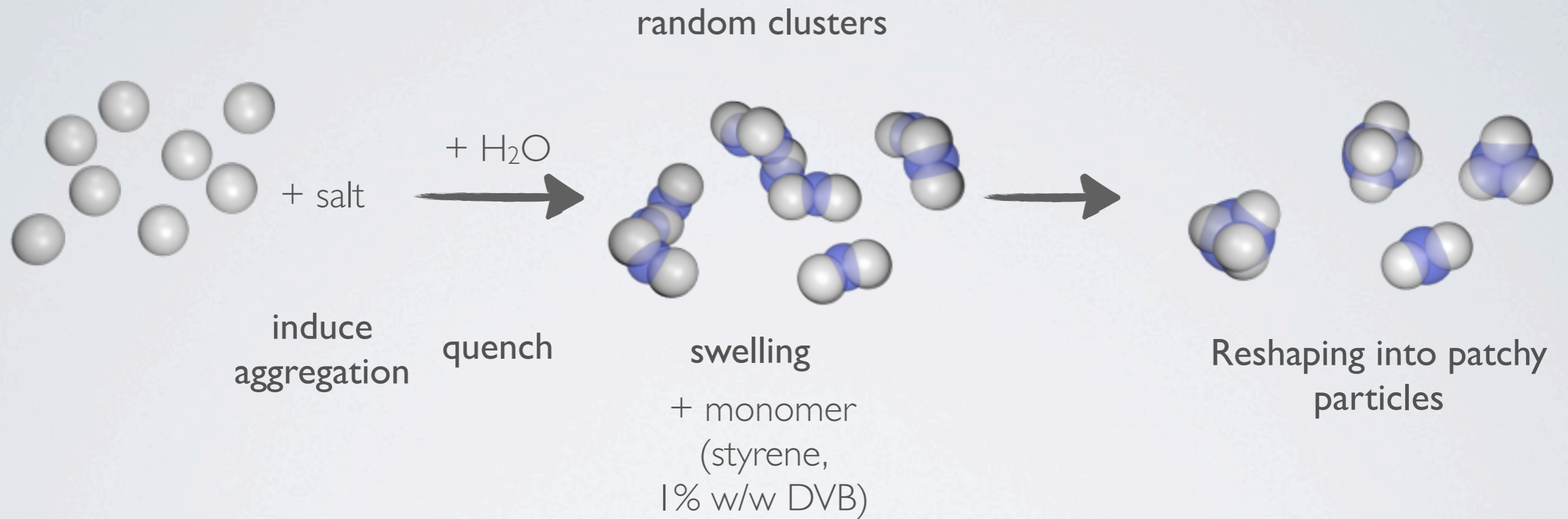
RESHAPING RANDOM COLLOIDAL CLUSTERS

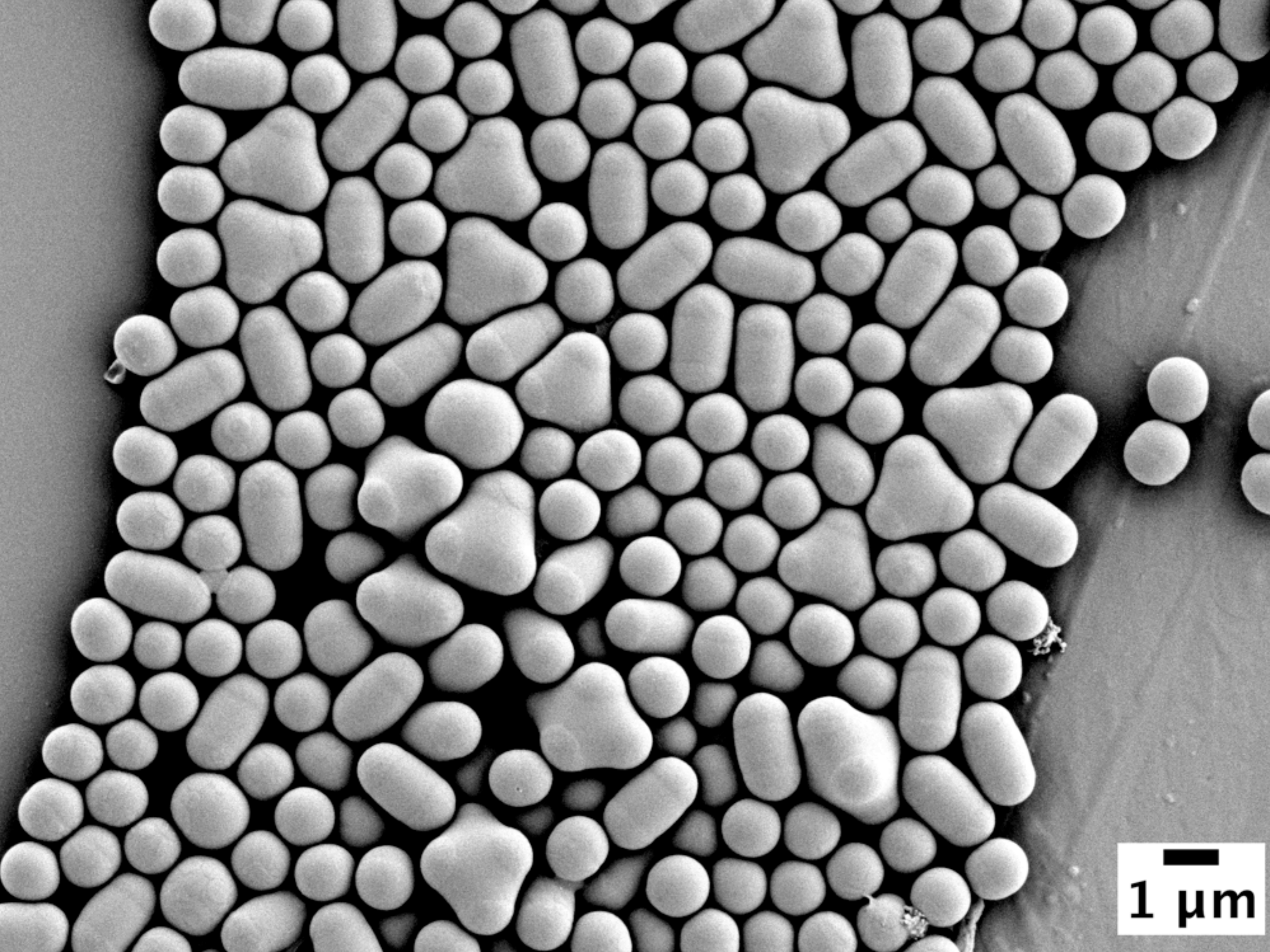


Charge stabilized
PS colloids

~ 1 μm diameter, 1% crosslink density,
purchased from Magsphere

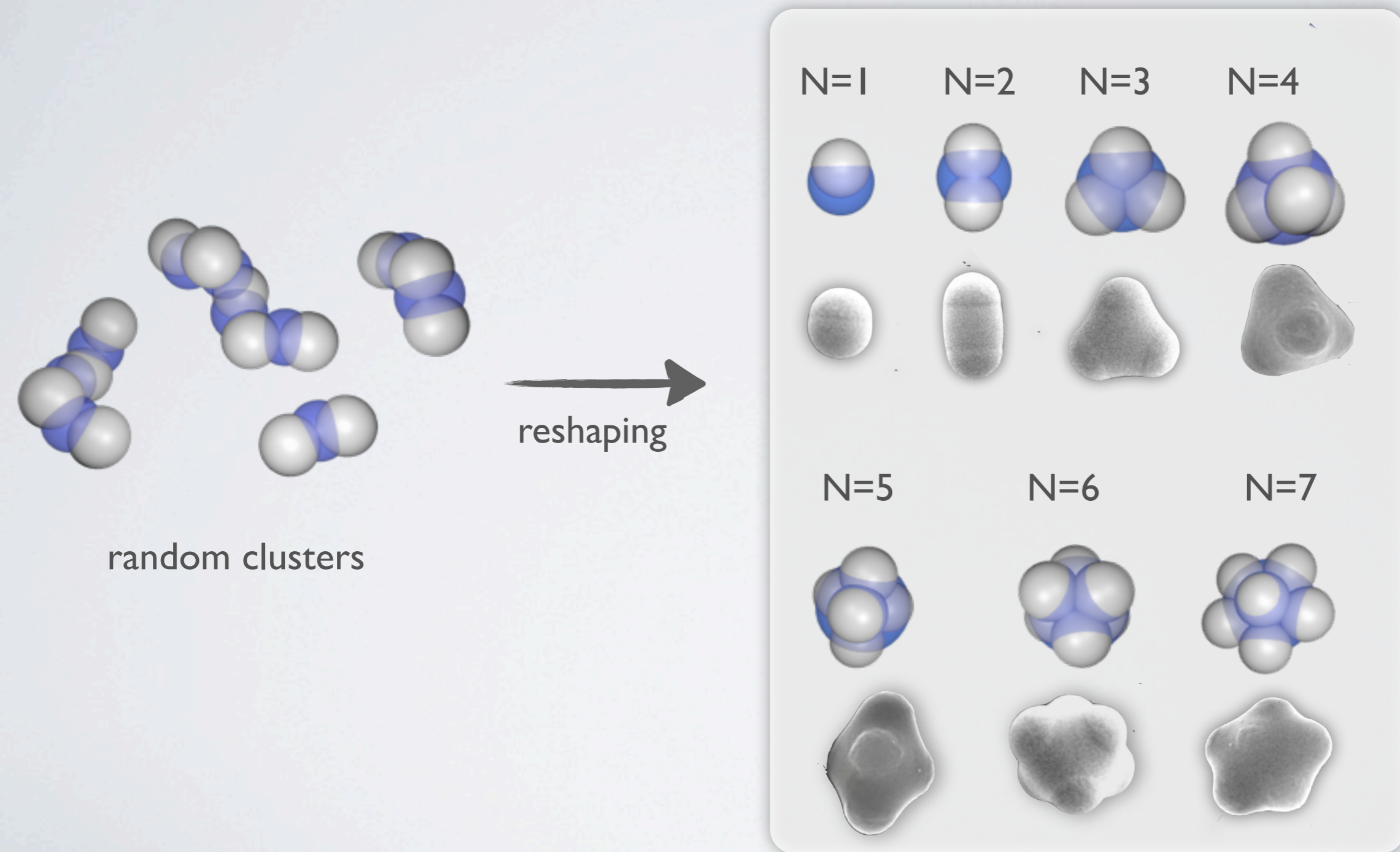
RESHAPING RANDOM COLLOIDAL CLUSTERS





—
1 μm

PARTICLE SWELLING RESHAPES THE RANDOM CLUSTERS INTO UNIFORM PATCHY PARTICLES

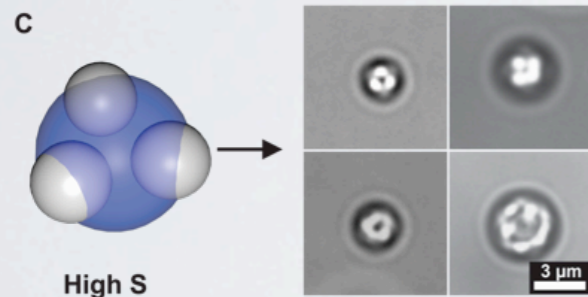


COALESCENCE DRIVEN RECONFIGURATION

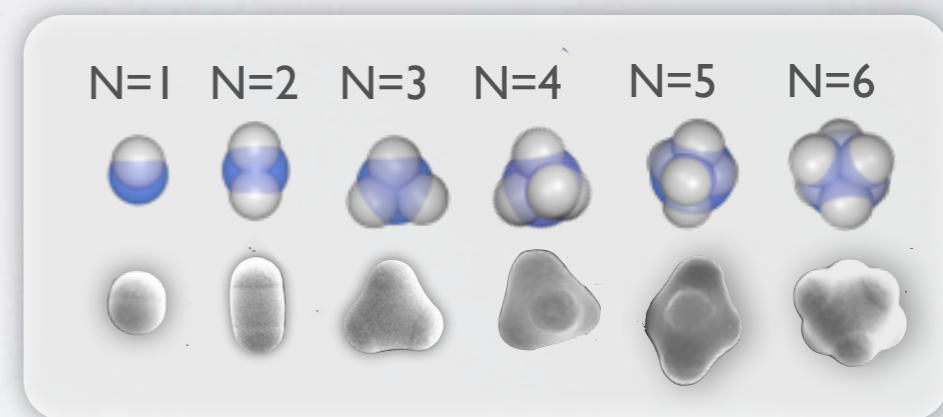
Liquid droplet coalescence drive rearrangement



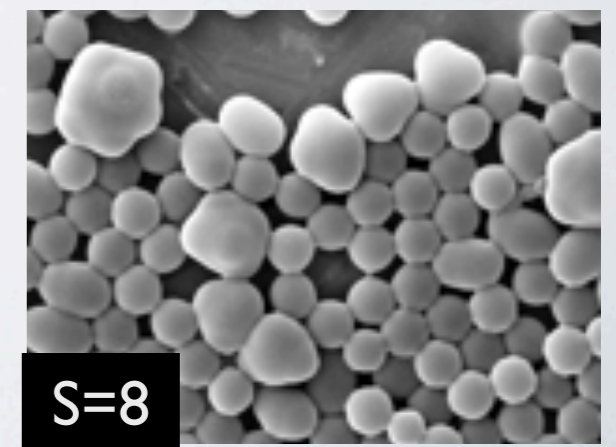
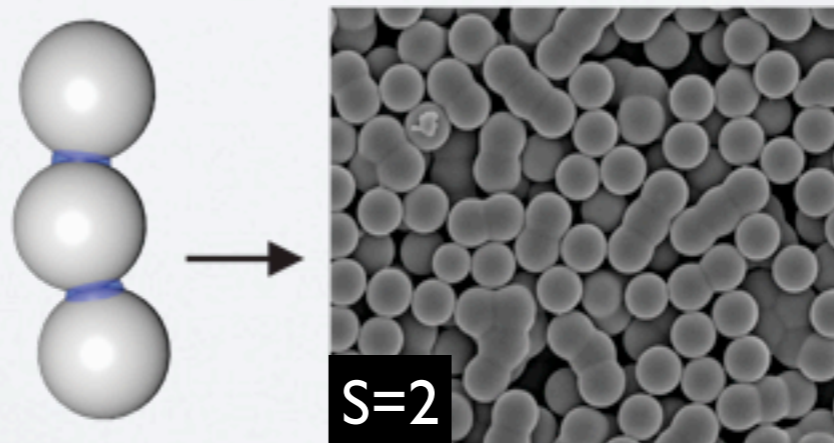
Liquid droplet confines the spheres



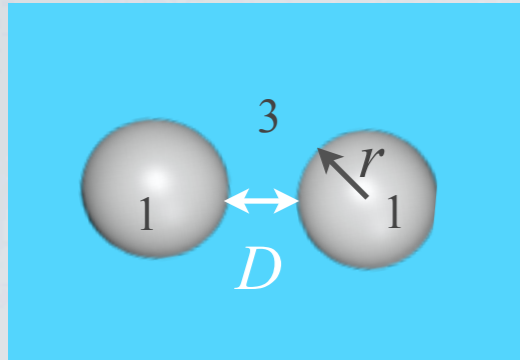
Cluster minimize the second moment of the mass distribution



Insufficient swelling
 → no / small liquid bridges
 → no reconfiguration!



WHAT ENABLES RECONFIGURATION?

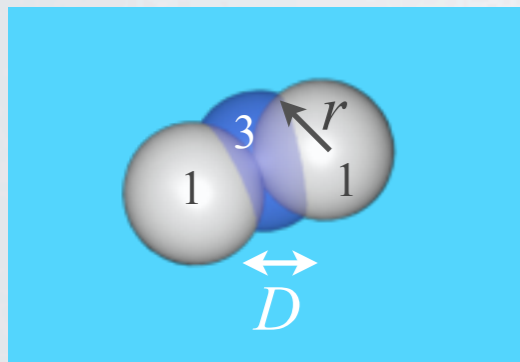


Van der Waals interaction energy $W(D) = -\frac{Ar}{12D}$
with the Hamaker constant A (Lifshitz theory)

$$A = \frac{3}{4}k_B T \left(\frac{\epsilon_1 - \epsilon_3}{\epsilon_1 + \epsilon_3} \right)^2 + \frac{3h\nu_e}{16\sqrt{2}} \frac{(n_1^2 - n_3^2)^2}{(n_1^2 + n_3^2)^{3/2}}$$

polystyrene spheres: $\epsilon_{PS} = 2.55$ $n_{PS} = 1.557$

polystyrene spheres in water: $\epsilon_w = 80$ $n_w = 1.333$ $A_{PS-w} = 1.5 \cdot 10^{-20} J$



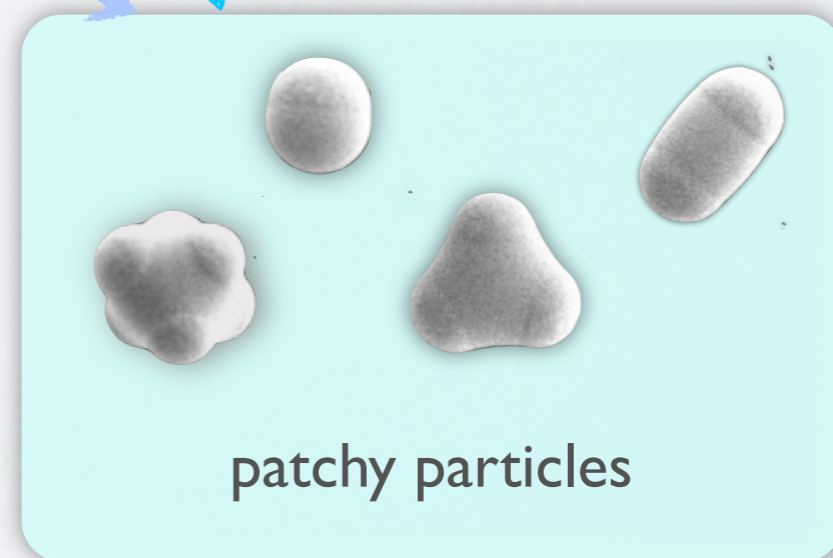
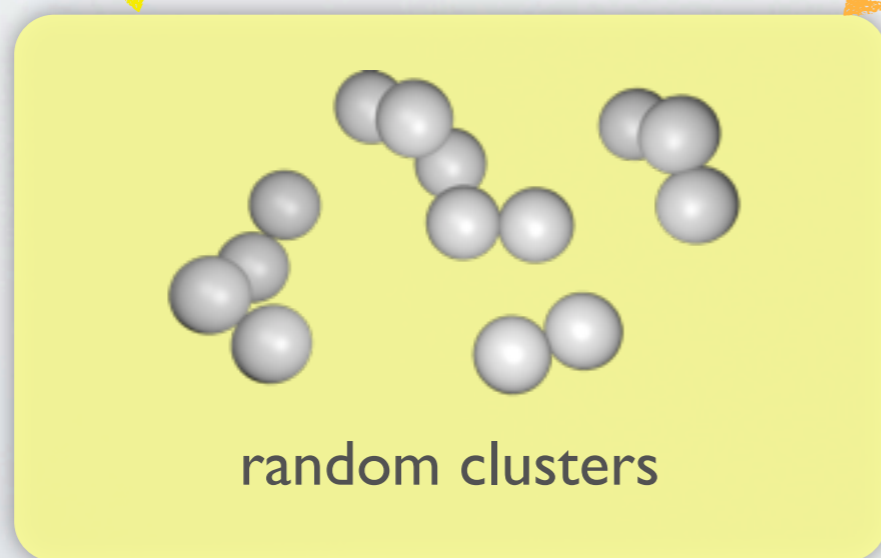
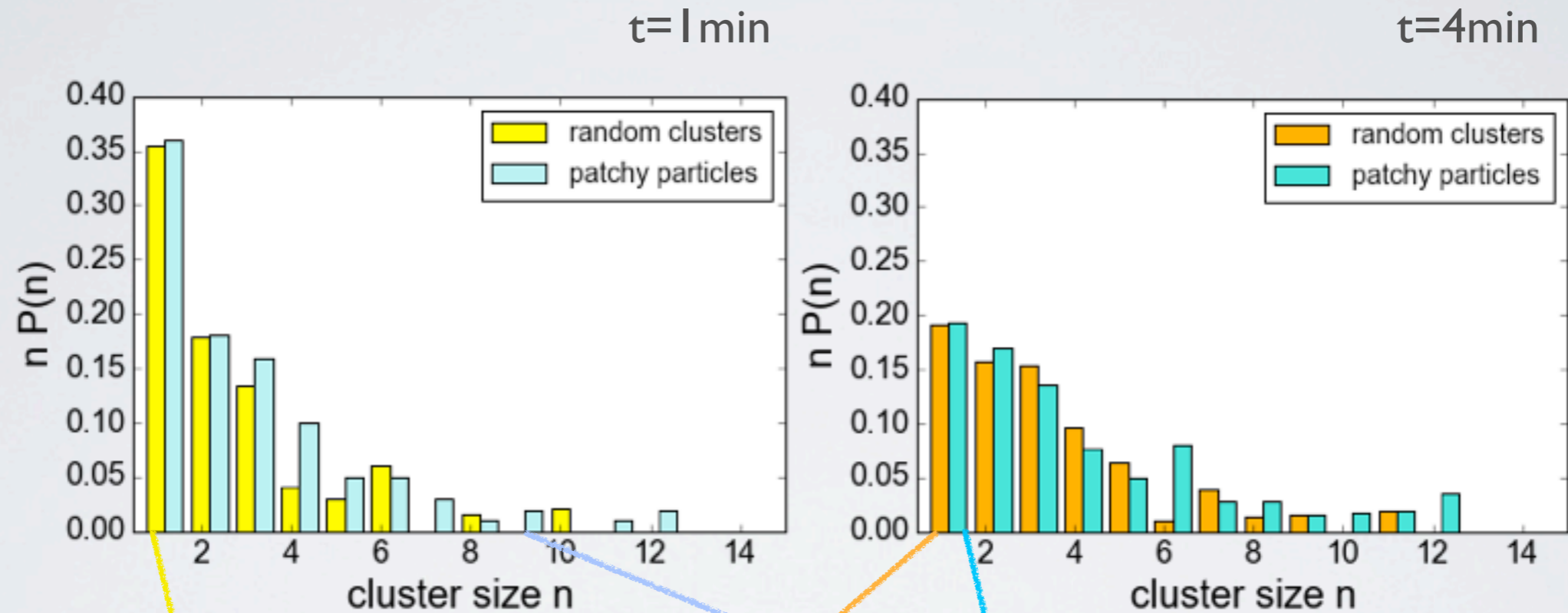
polystyrene spheres in styrene:

$\epsilon_{st} = 2.8$ $n_{st} = 1.448$

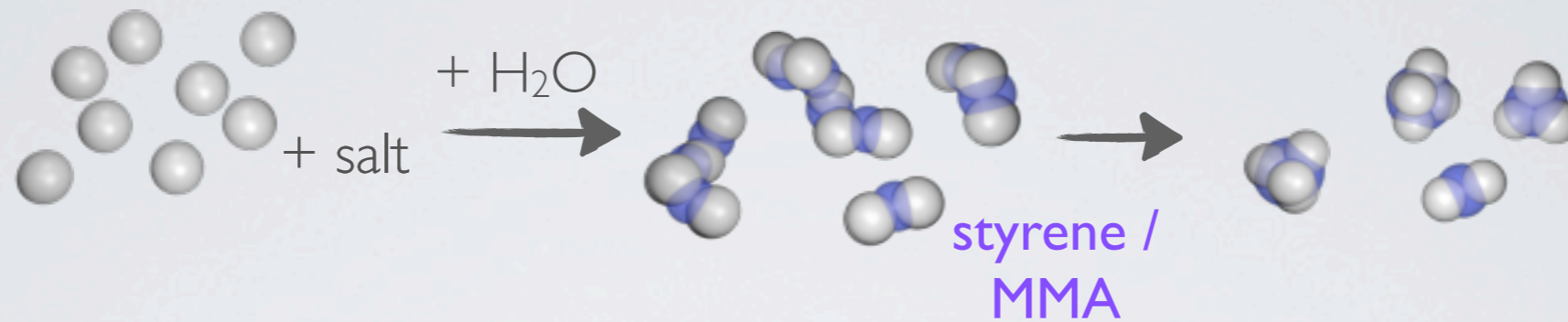
$A_{PS-st} = 5.3 \cdot 10^{-23} J$

Significant reduction of van der Waals attraction due to liquid bridges!

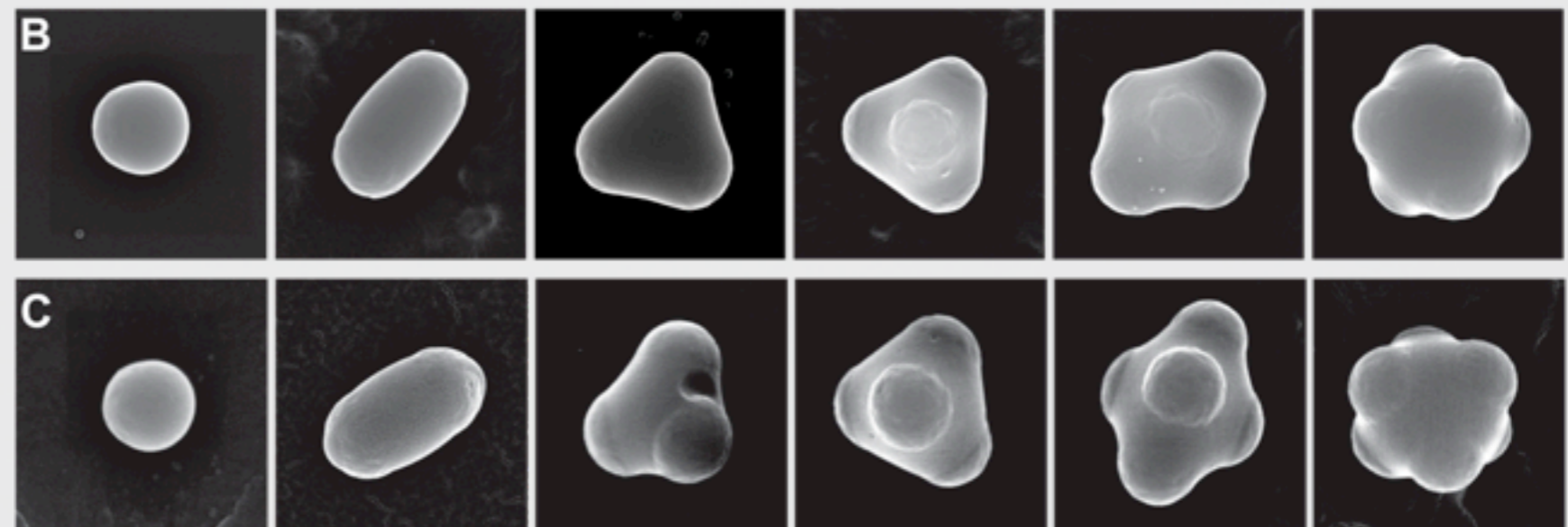
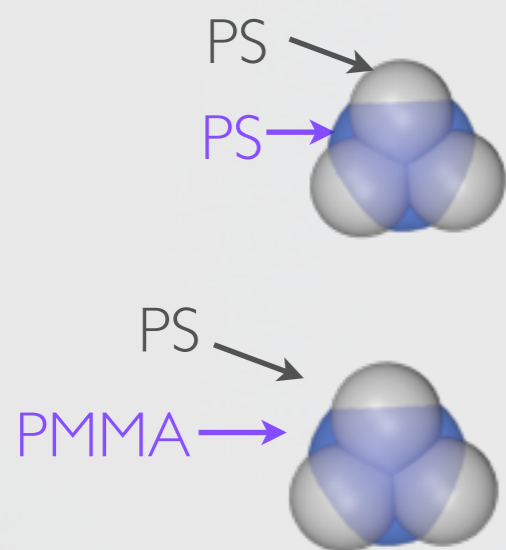
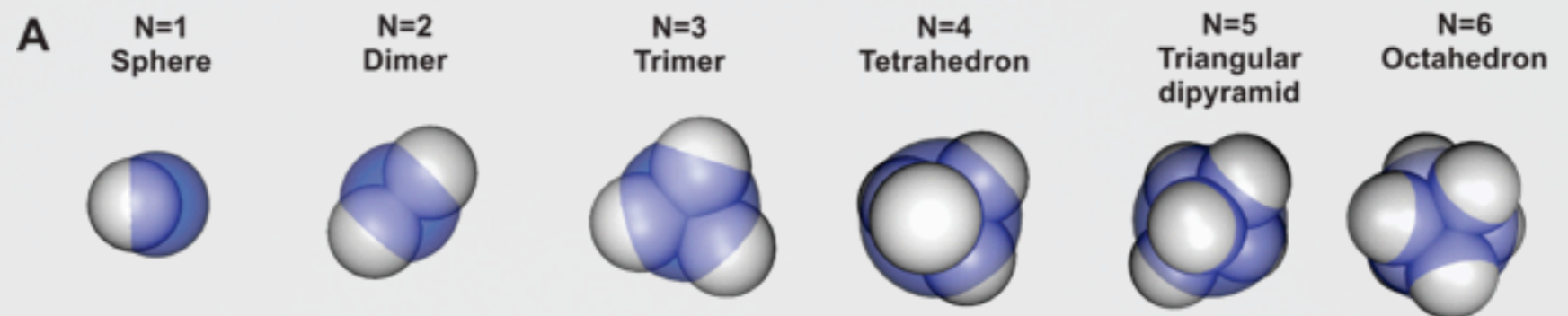
SIZE DISTRIBUTION OF RANDOM CLUSTERS AND PATCHY PARTICLES IS EQUAL



COMPOSITE PS / PMMA COLLOIDAL MOLECULES

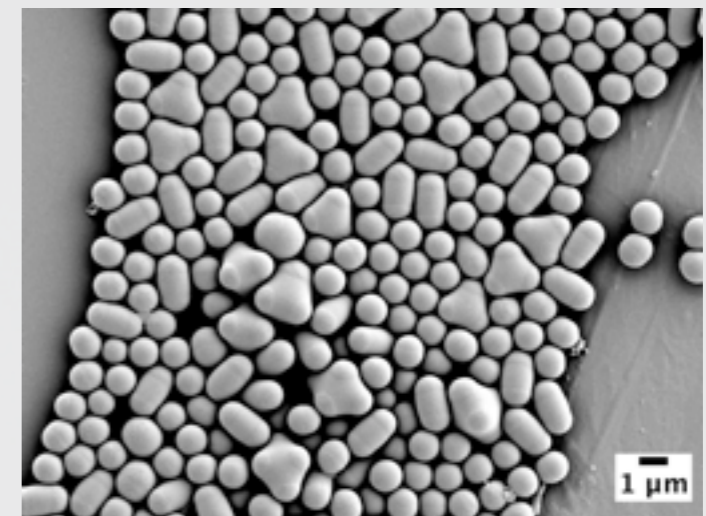
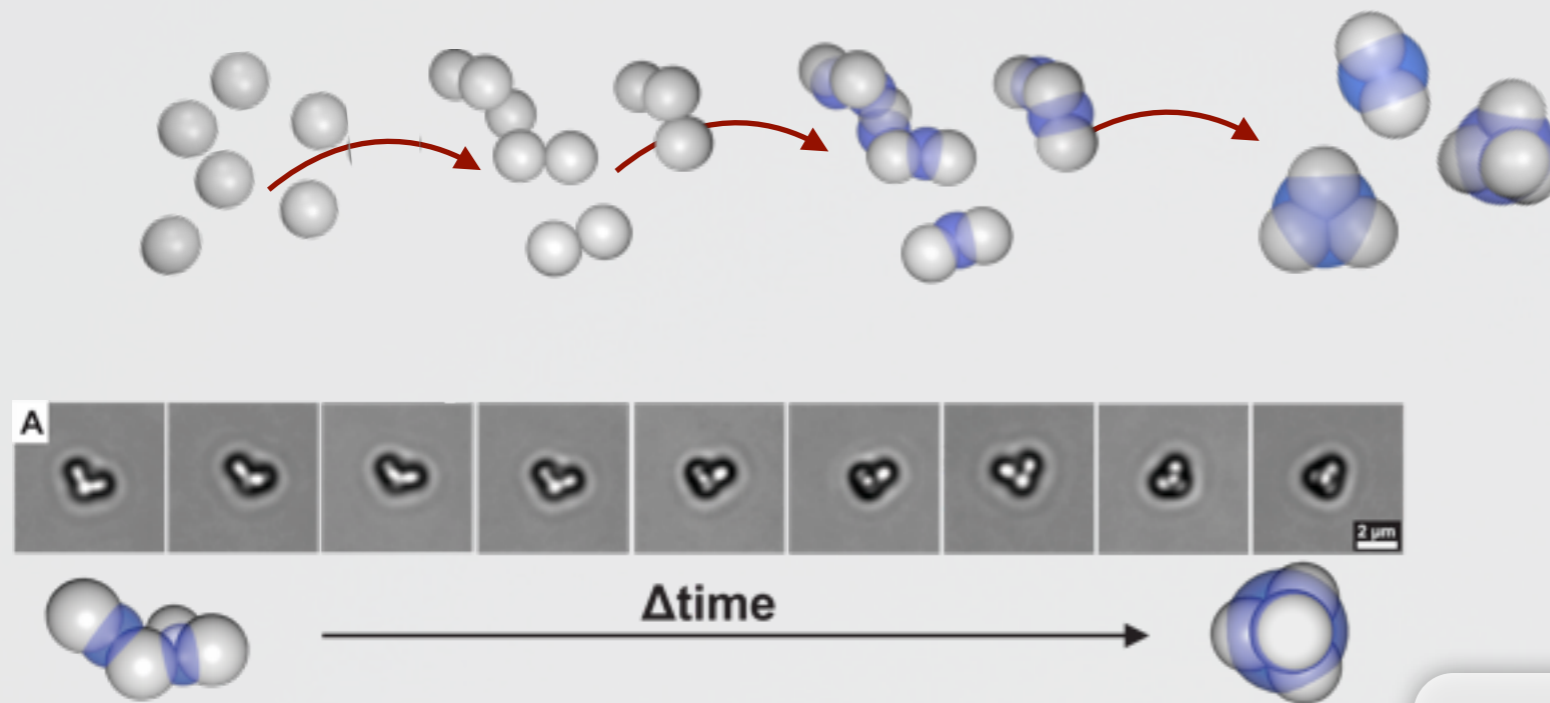


Patchy particles



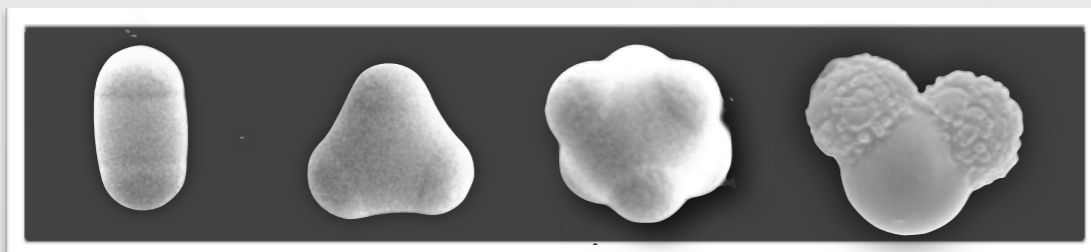
SUMMARY - RECYCLING COLLOIDAL AGGREGATES INTO PATCHY PARTICLES

Reorganization of random clusters of spheres

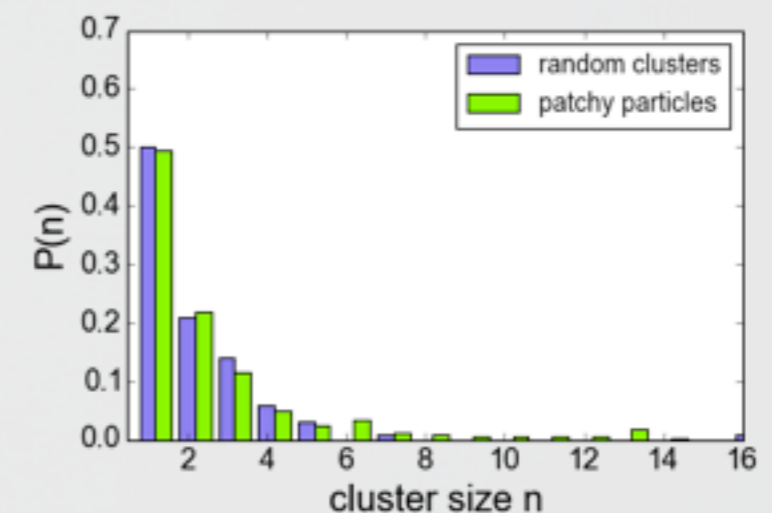


Meester, Verweij, vd Wel, Kraft (ACSNano 2016)

Variety of complex patchy particles



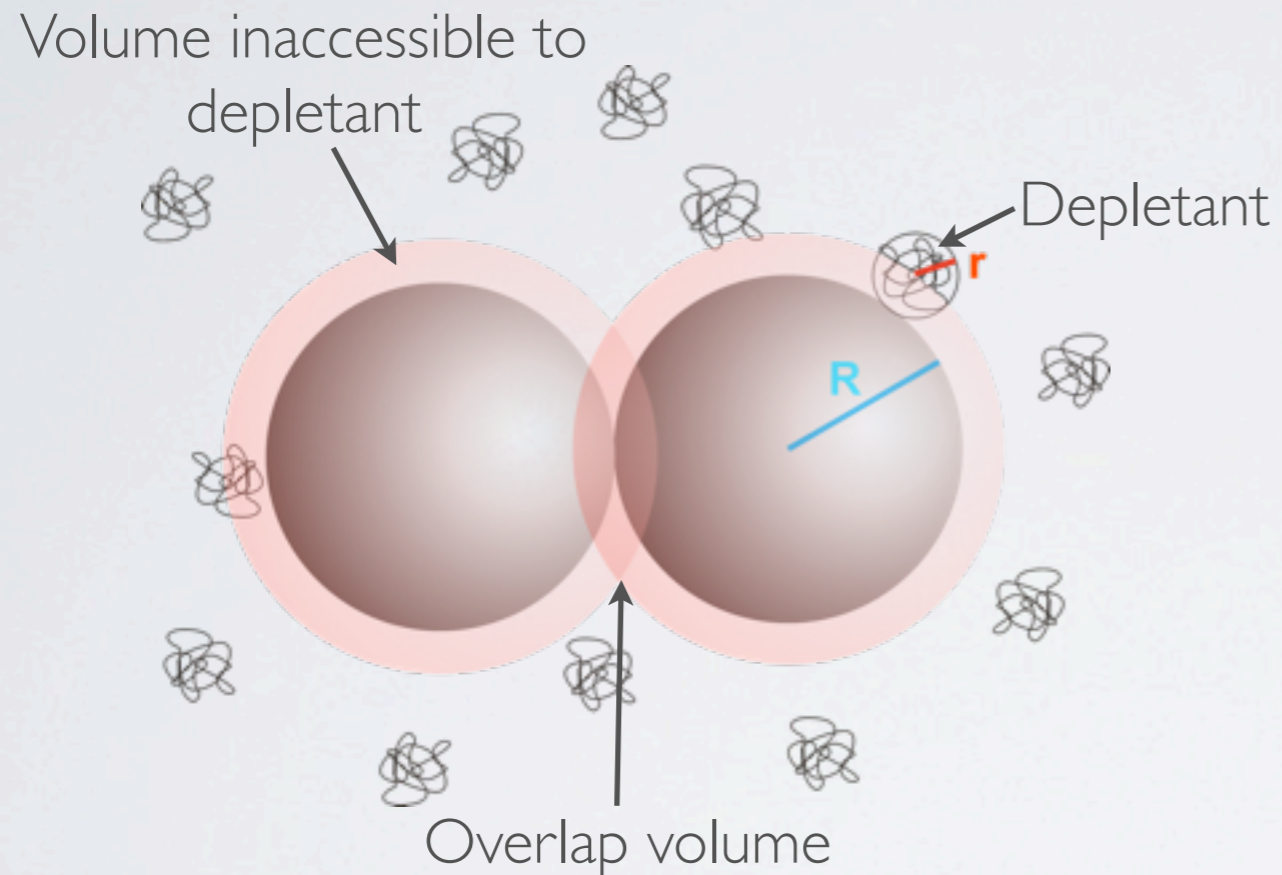
Control over size distribution



The background of the slide is a close-up photograph of several orange segments, showing their characteristic bumpy, porous texture. The segments are arranged in a somewhat circular pattern, with some in the foreground and others slightly behind, creating a sense of depth. The lighting is soft, highlighting the natural colors and textures of the fruit.

SELF-ASSEMBLY OF PATCHY PARTICLES
- TUNING DEPLETION INTERACTIONS BY
SURFACE ROUGHNESS

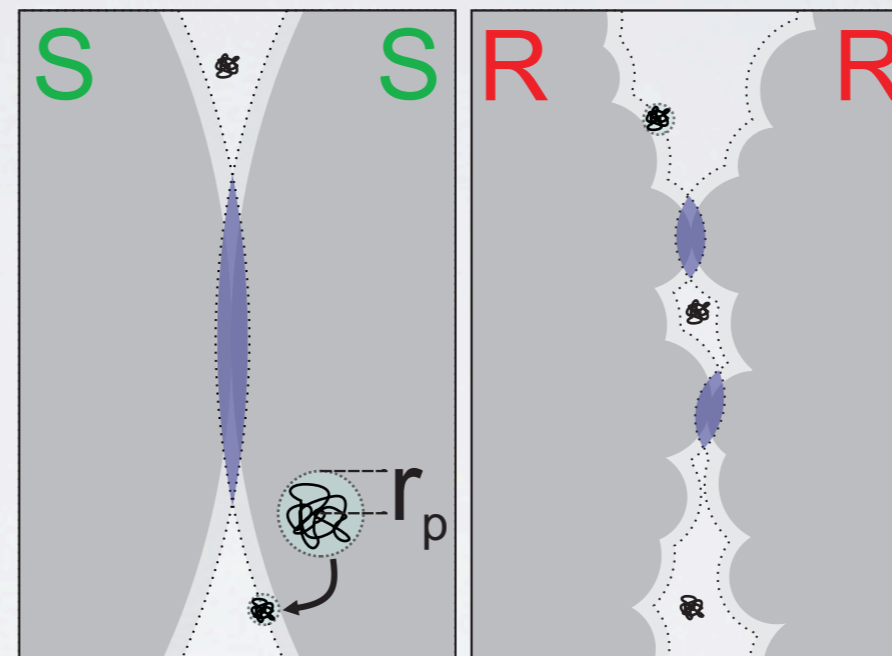
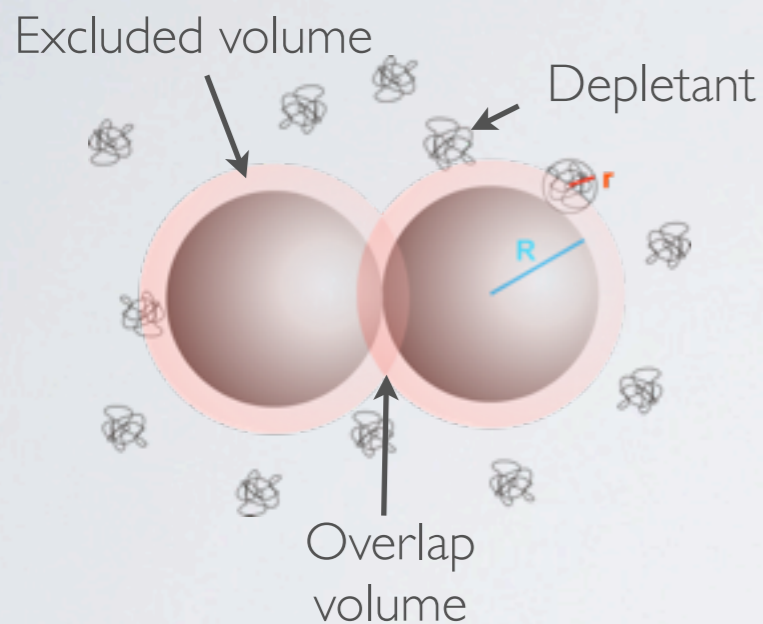
THE DEPLETION INTERACTION



Effective attraction between larger colloid!

$$\begin{aligned} u &= -\Pi V_{overlap} \\ &= -k_B T n_{depl} V_{overlap} \end{aligned}$$

TUNING DEPLETION INTERACTIONS THROUGH THE OVERLAP VOLUME



$$\begin{aligned}
 u &= -\Pi V_{overlap} \\
 &= -k_B T n_{depl} V_{overlap}
 \end{aligned}$$

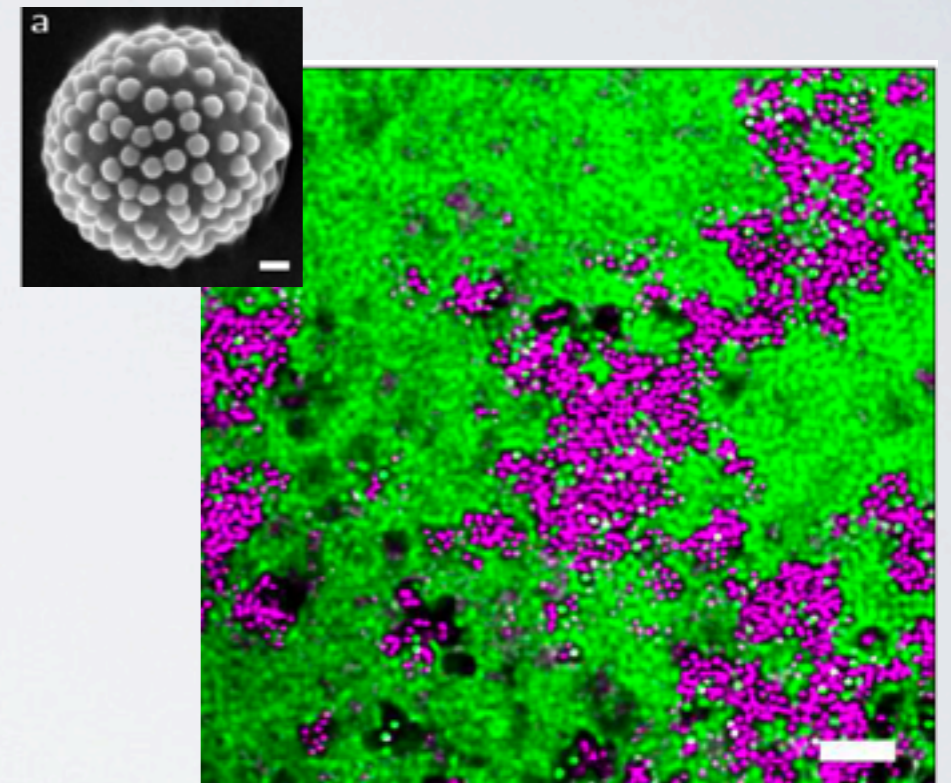
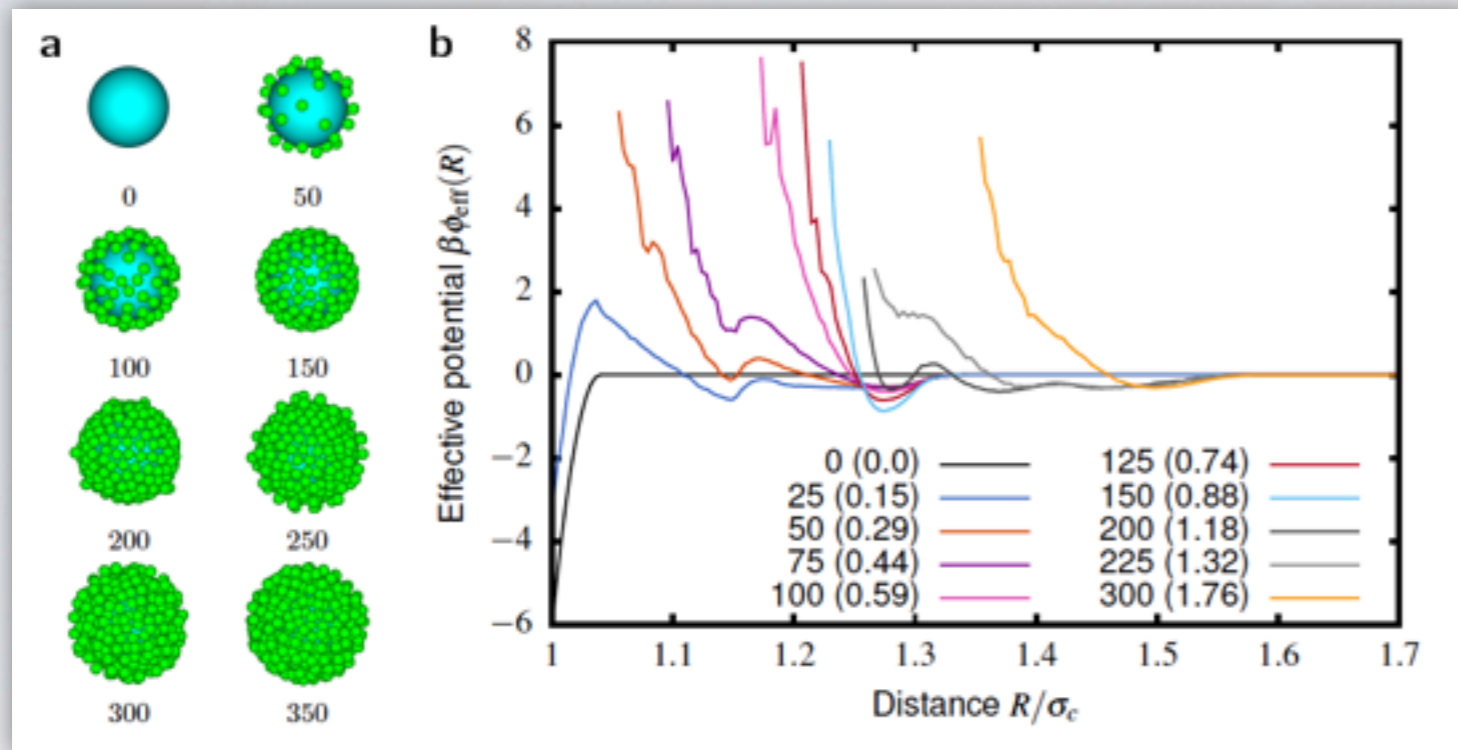
Roughness significant reduces the overlap volume, and thus the depletion attraction.

see also: K. Zhao, T. Mason, PRL 99(2008)

Asakura, Oosawa, J. Chem. Phys. (1954)

Vrij Pure & Appl Chem. (1976)

DEPLETION POTENTIAL FOR ROUGH AND SMOOTH SPHERES



averaged over 60 frames

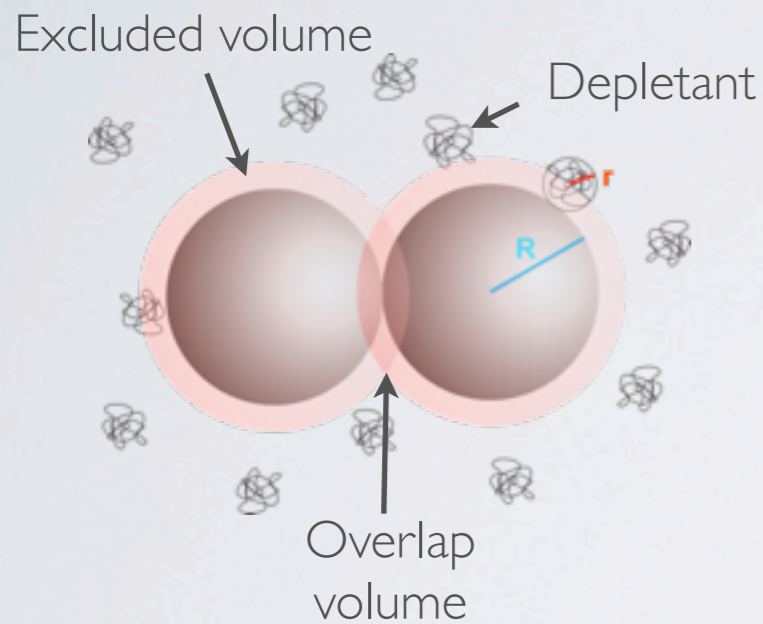
purple: smooth particles
green: rough particles

Simulations by Michiel Hermes

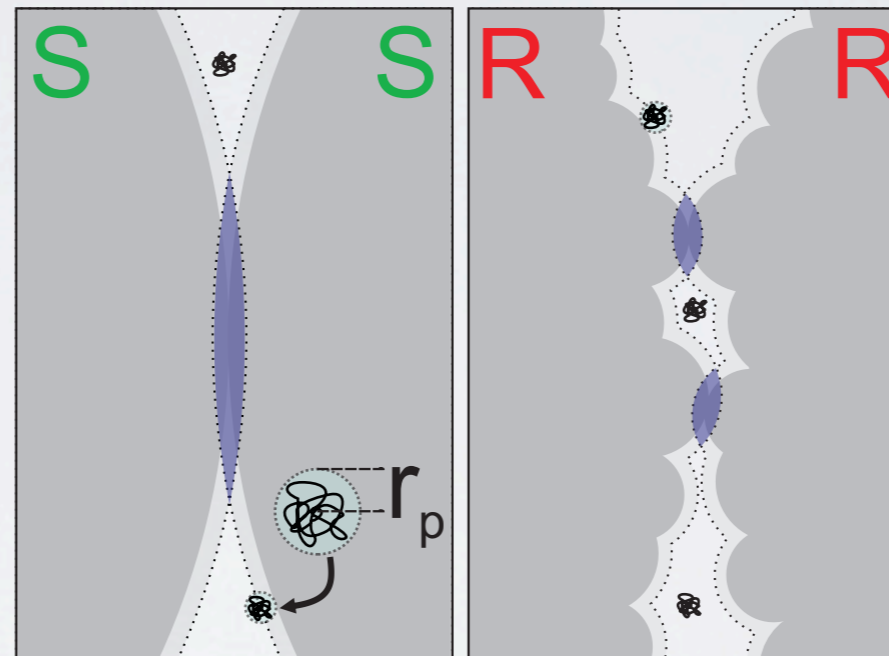
volume fraction=0.16

$\sigma_p=0.04$ σ_c

TUNING DEPLETION INTERACTIONS THROUGH THE OVERLAP VOLUME

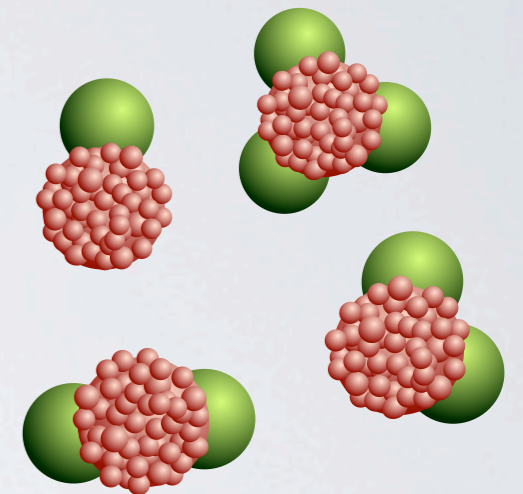


$$\begin{aligned}
 u &= -\Pi V_{overlap} \\
 &= -k_B T n_{depl} V_{overlap}
 \end{aligned}$$



Roughness significant reduces the overlap volume, and thus the depletion attraction.

**smooth =
large overlap volume =
attractive**



**rough =
small overlap volume =
non-attractive**

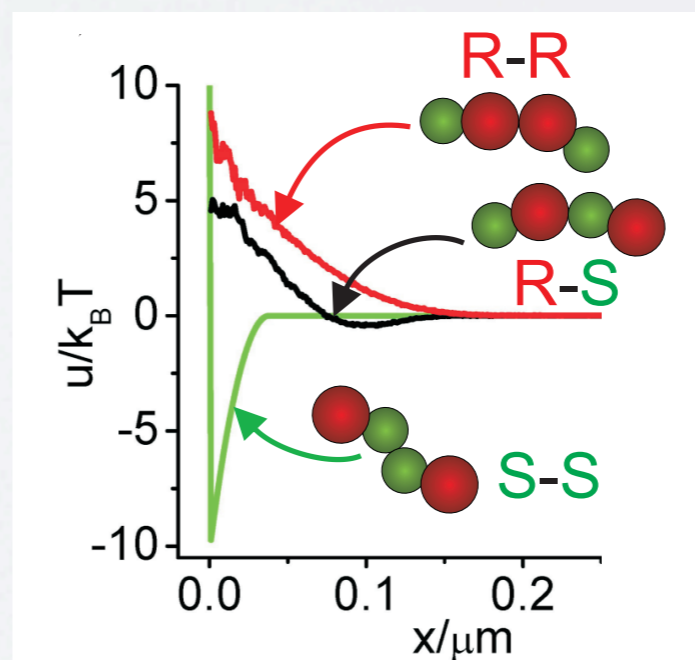
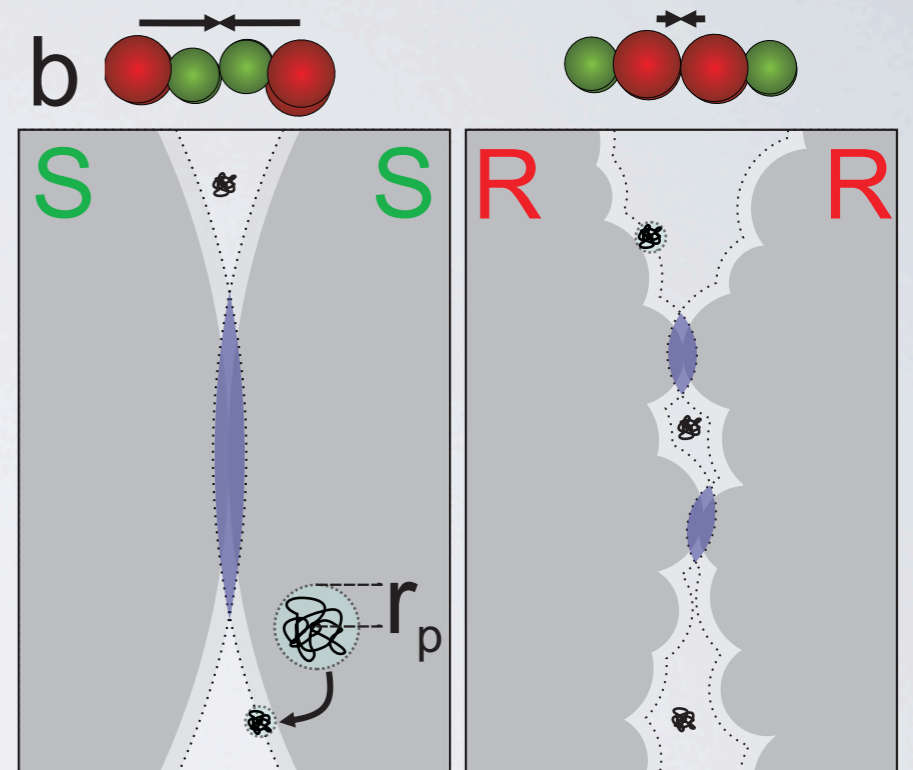
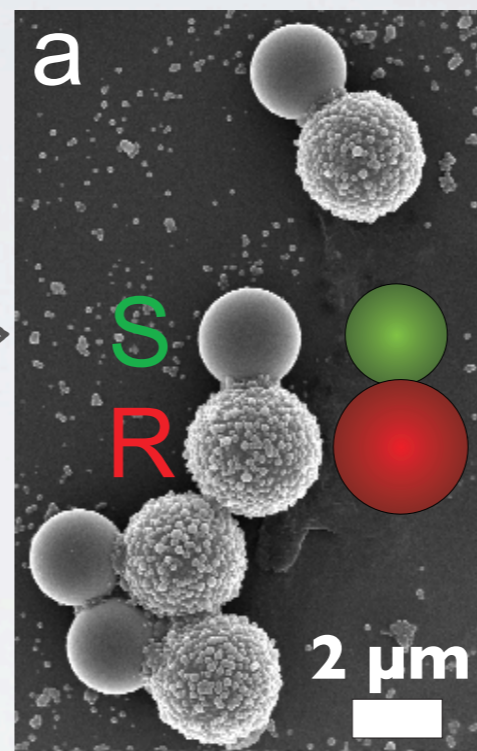
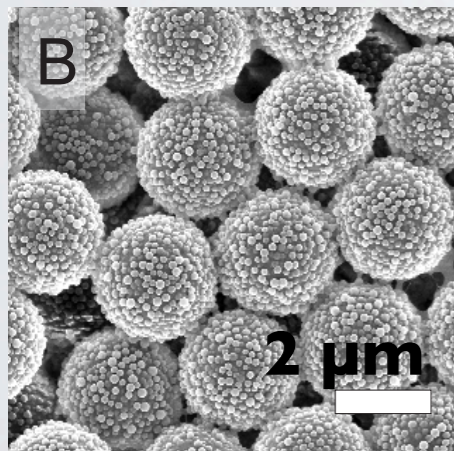
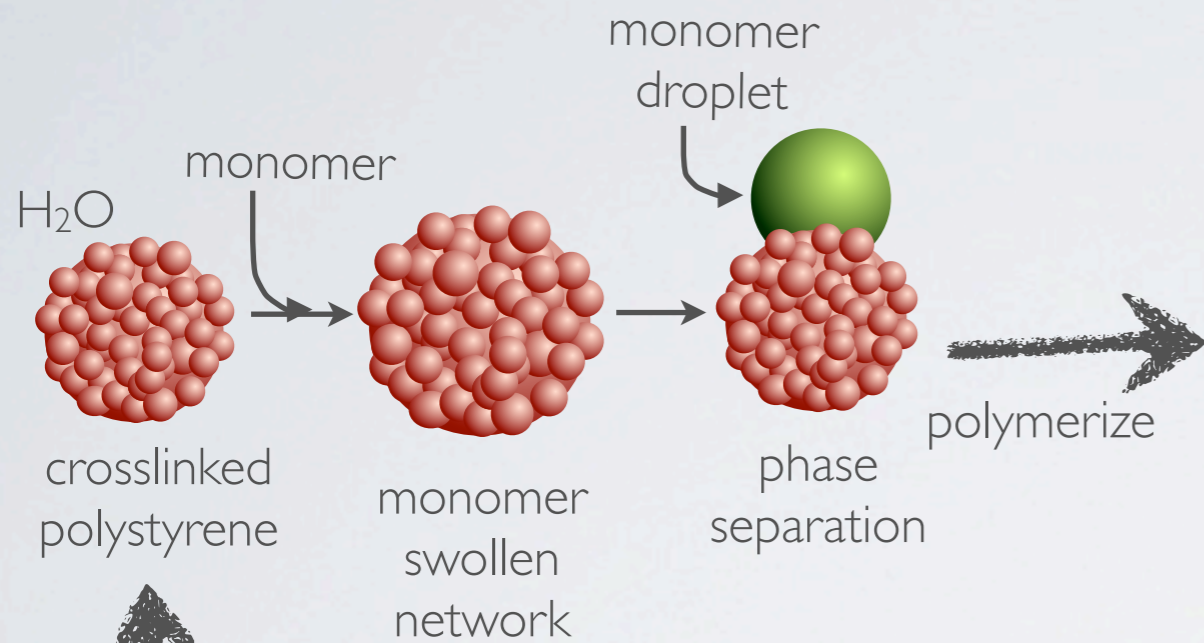
Patchy Particles

see also: K. Zhao, T. Mason, PRL 99(2008)

Asakura, Oosawa, J. Chem. Phys. (1954)

Vrij Pure & Appl Chem. (1976)

SYNTHESIS OF ROUGH-SMOOTH COLLOIDS

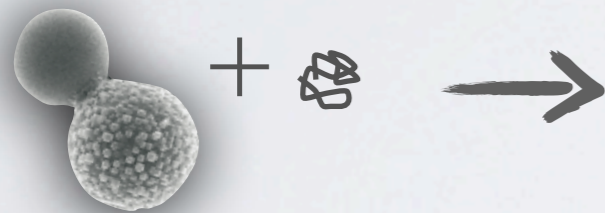


Significant attraction between smooth surfaces only

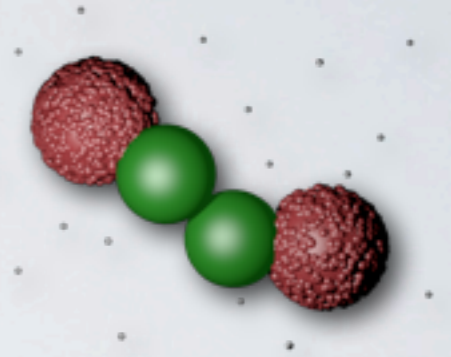
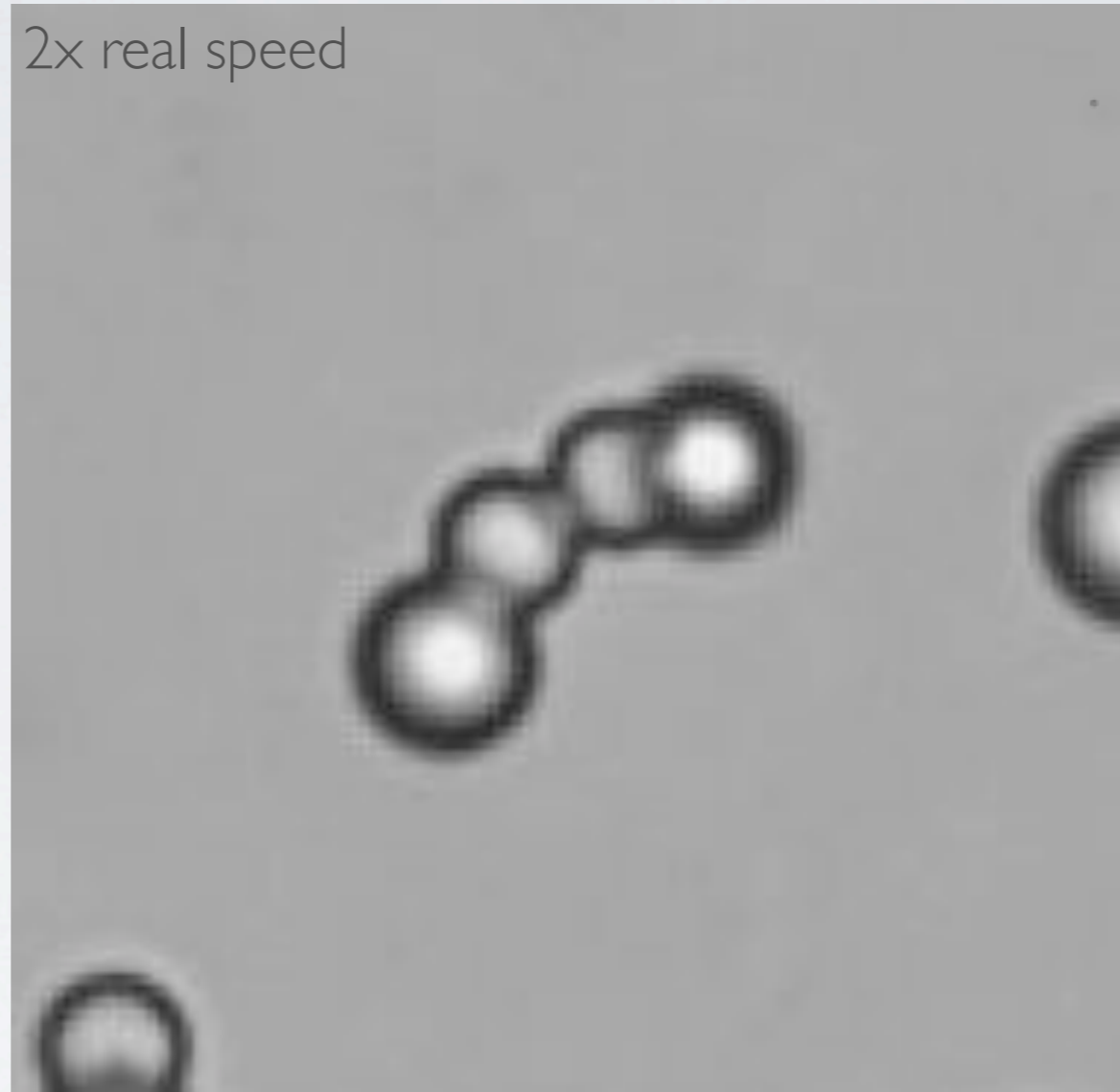
Potential calculated by M. Hermes and M. Dijkstra

BOND FORMATION AND BREAKING

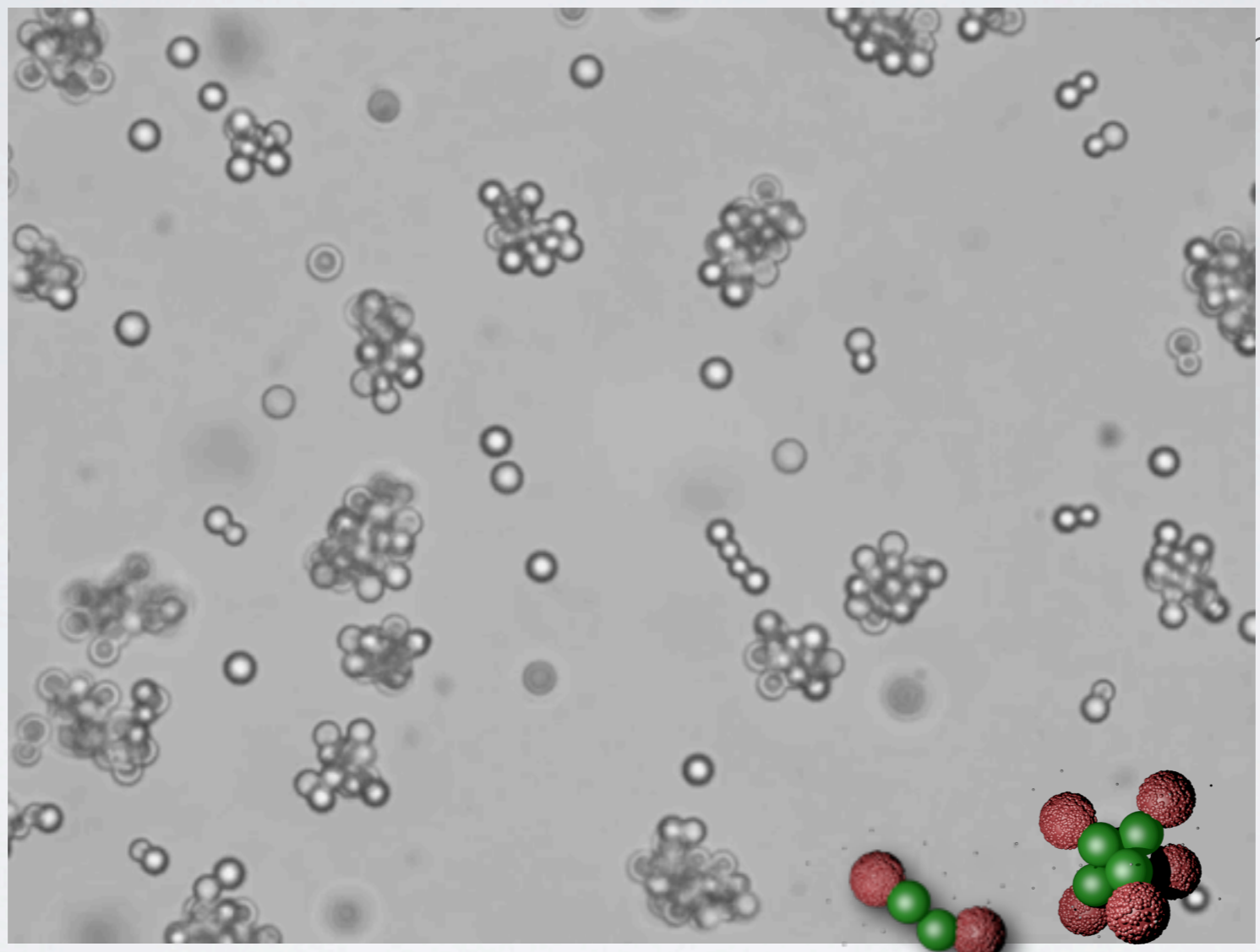
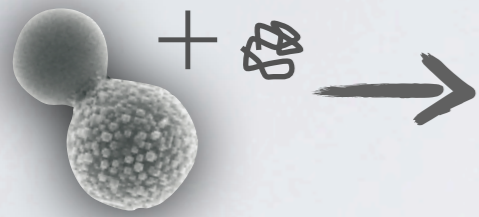
Roughness anisotropic colloids +
Dextran as depletant ($d=38\text{nm}$)



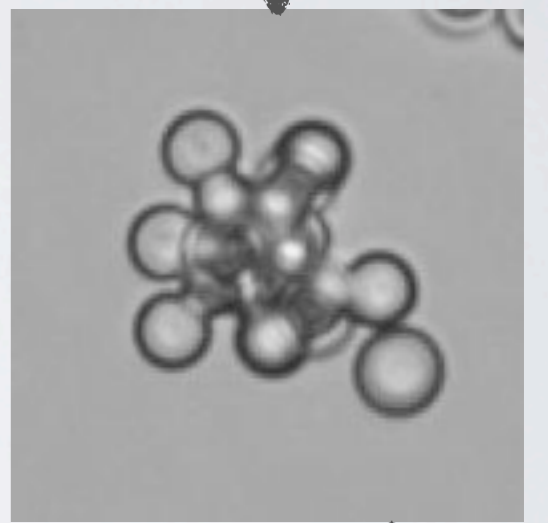
- Site-specific binding
- Flexible bonds



COLLOIDAL MICELLES



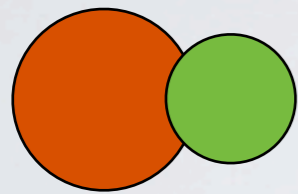
Cluster of colloids with smooth and attractive sides inside



Rough sphere (not attractive)

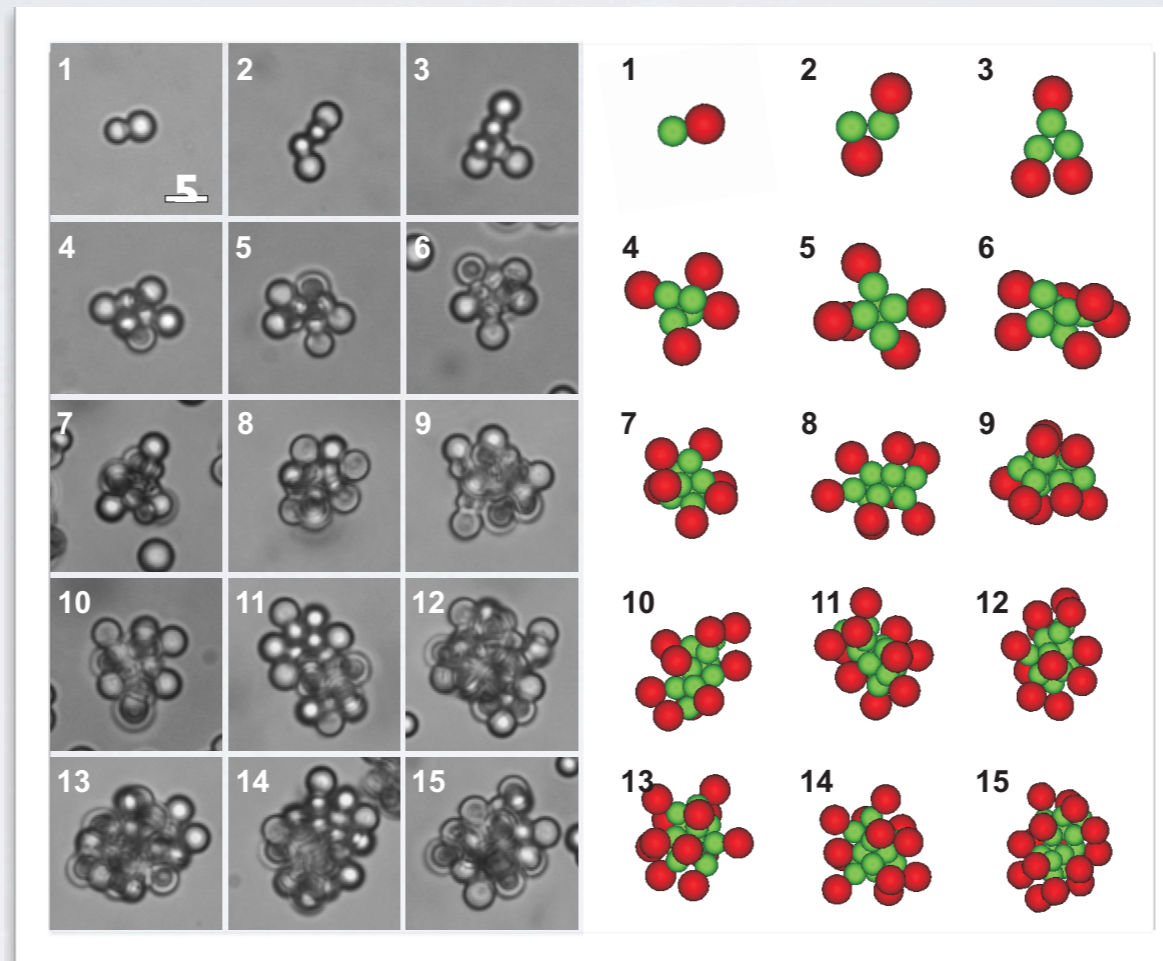
EXPERIMENT & SIMULATIONS

Monte-Carlo Simulations by R. Ni & M. Dijkstra

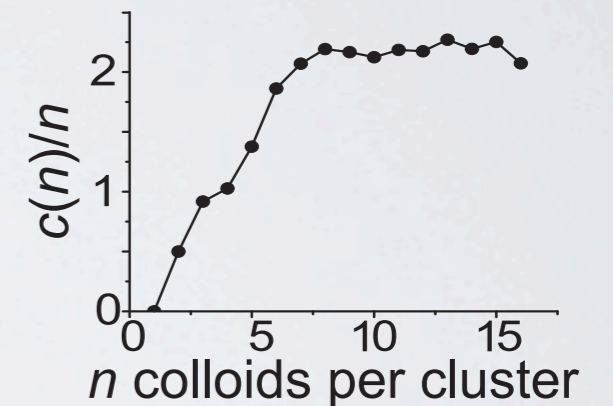


hard sphere
repulsion

Asakura-
Oosawa-Vrij
potential



□ Average number of bonds saturates



□ Interaction potential
-9.85kT

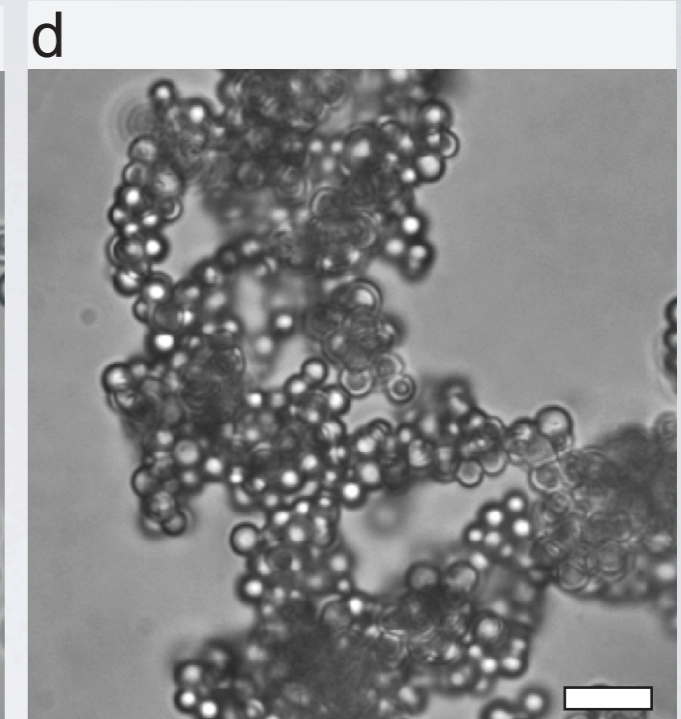
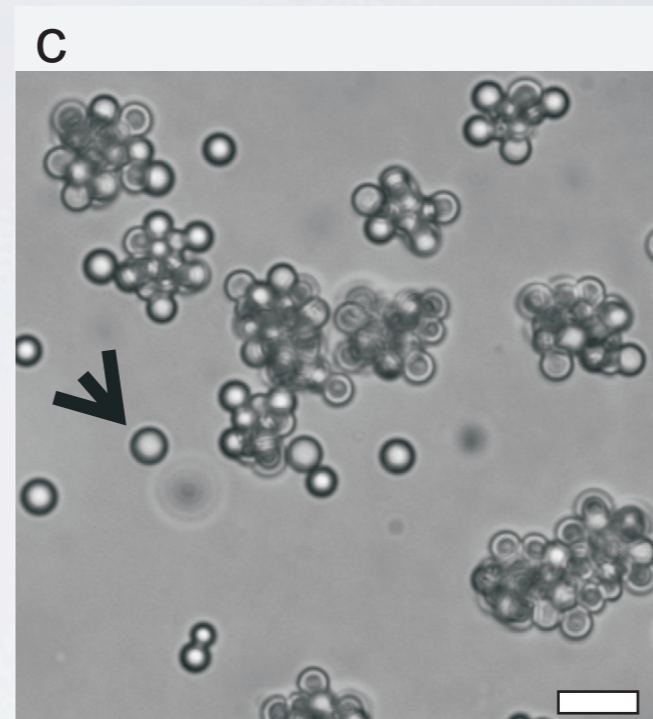
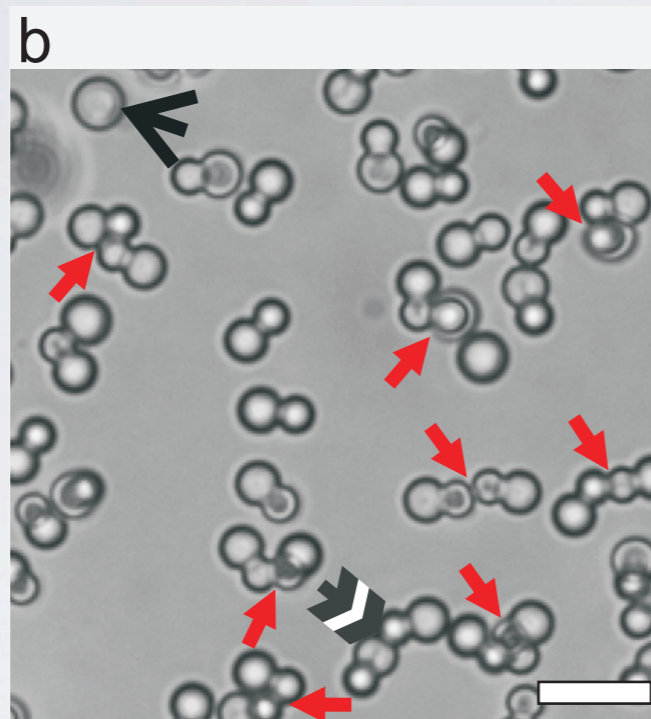
INCREASING INTERACTION STRENGTH

$\varphi = 0.32 \varphi_0$

$\varphi = 0.35 \varphi_0$

$\varphi = 0.40 \varphi_0$

$\varphi = 0.42 \varphi_0$

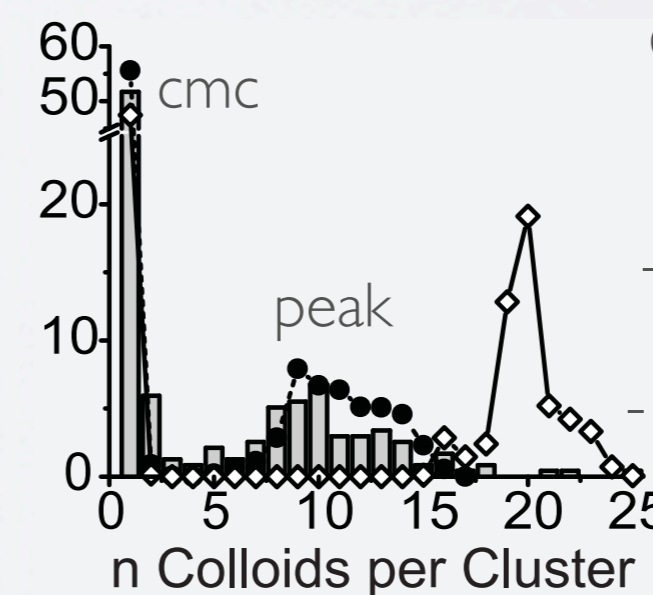
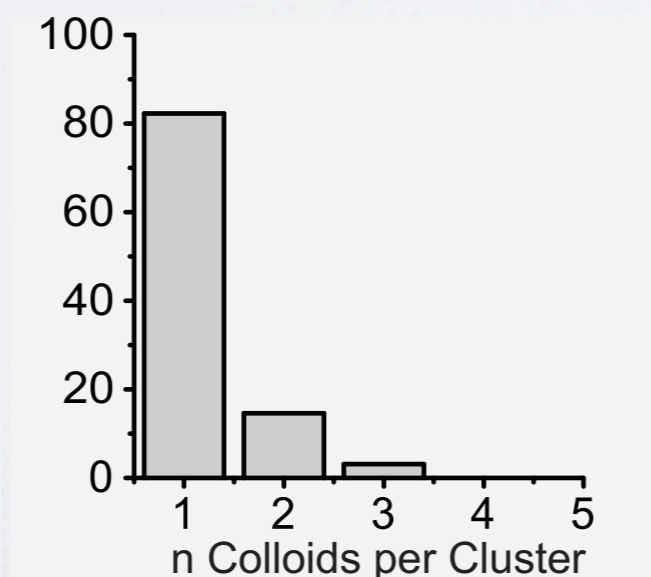
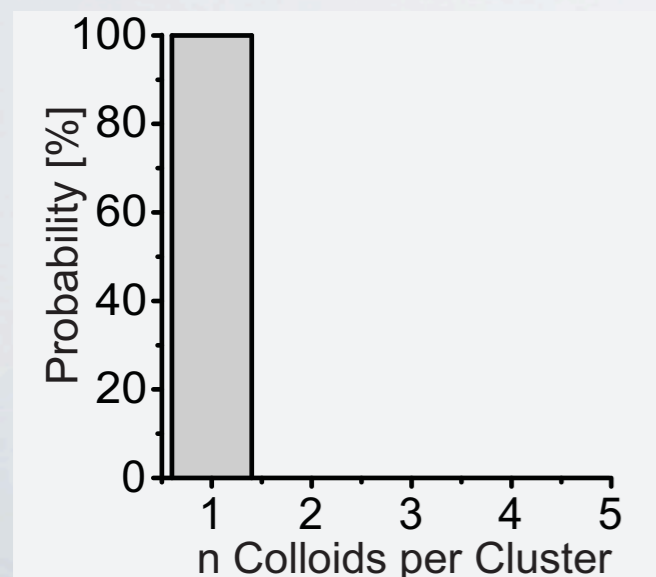


free particles

single bonds

cluster phase

gelation, site-specificity lost



φ_0 = overlap concentration

Experiment

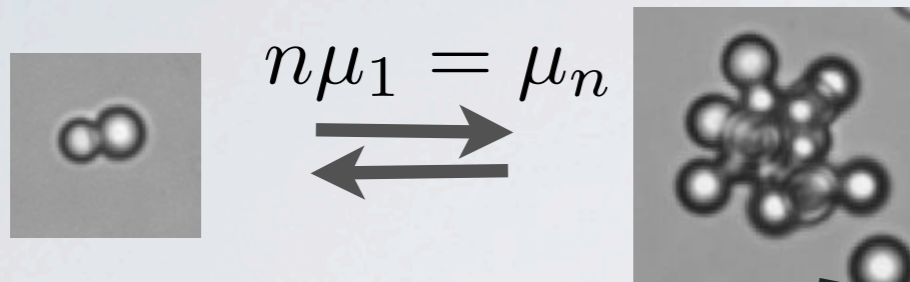
MC simulations
(R. Ni & M. Dijkstra)

Free Energy calculations
(F. Smallenburg & M. Dijkstra)

Kraft et al. PNAS (2012)

CRITICAL MICELLE CONCENTRATION

Analyze equilibria between monomers and clusters



Critical micelle concentration

$$\Phi_{th}^{cmc} = \frac{V_c}{\xi^3} \exp \left[\frac{\beta c(n^*)u}{(n^* - 1)} \right]$$

volume of a colloid $\rightarrow V_c$
 interaction range $\rightarrow \xi$
 $\beta c(n^*)u$ \rightarrow # bonds in a cluster of size n^* \times bond energy
 $\frac{\beta c(n^*)u}{(n^* - 1)}$ \rightarrow \cong bond energy per particle in a cluster

Experiments

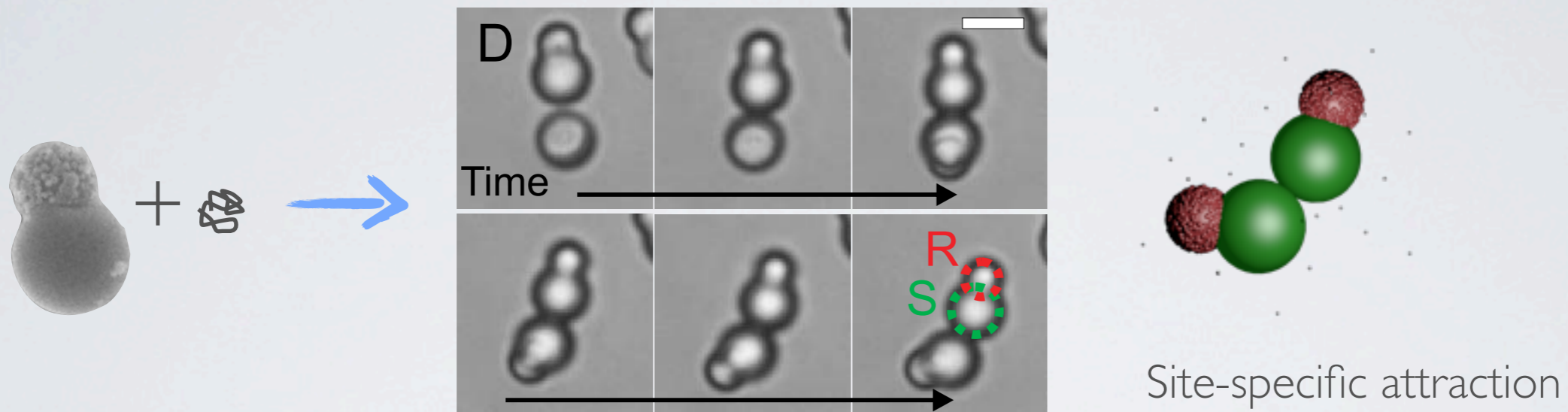
$$\Phi_{exp}^{cmc} = 3.1 \cdot 10^{-5}$$

Theoretical prediction

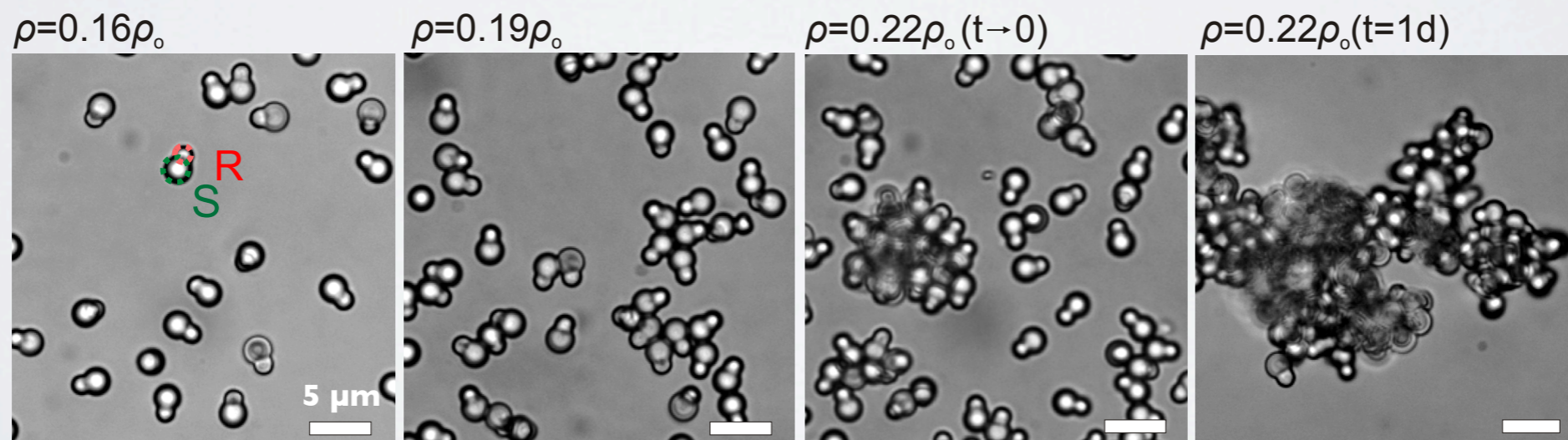
$$\Phi_{MC}^{cmc} = 1.3 \cdot 10^{-5}$$

Interaction range & energy determine the critical micelle concentration

INFLUENCE OF PATCH SIZE



□ Increasing attraction strength



free particles

single bonds

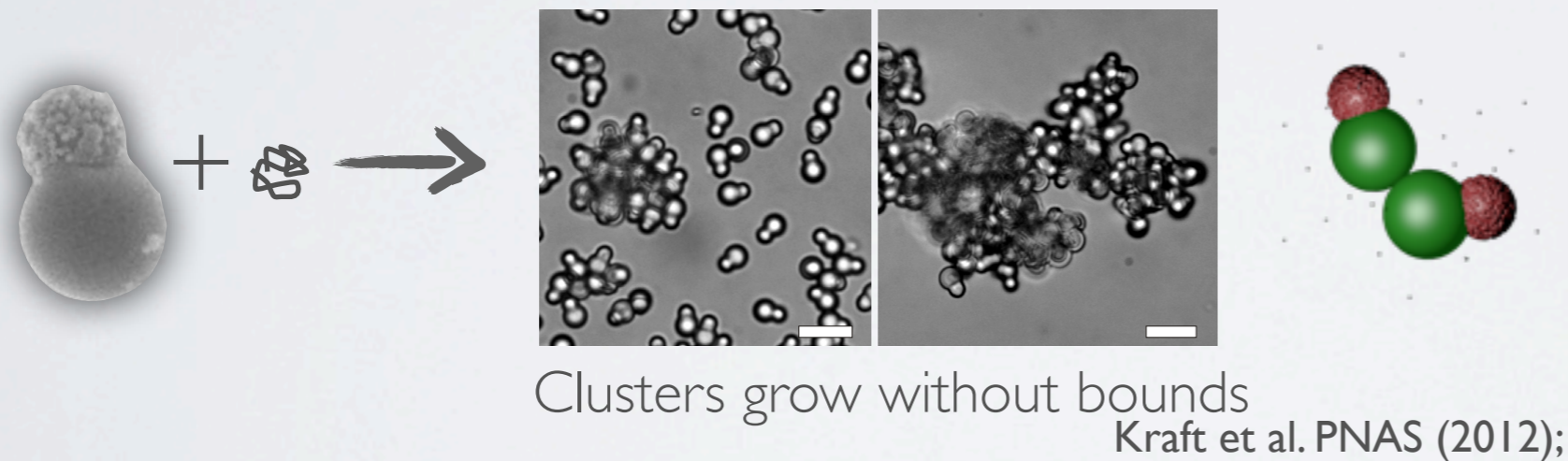
clusters grow without bounds

$\varphi_0 = \text{overlap concentration}$

STERIC CONSTRAINTS DETERMINE SIZE AND GEOMETRY OF THE STRUCTURE

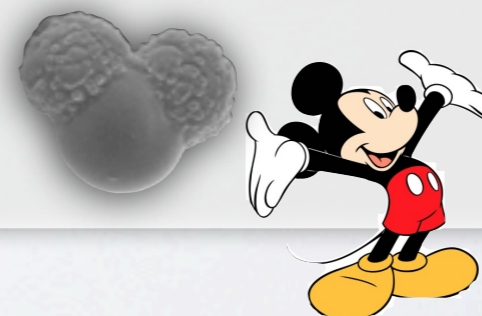
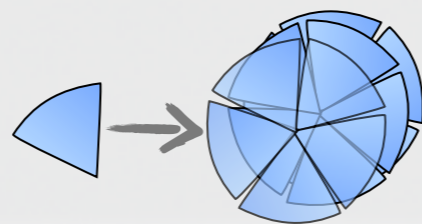


Rough ends sterically constrain assembly size in 3D.



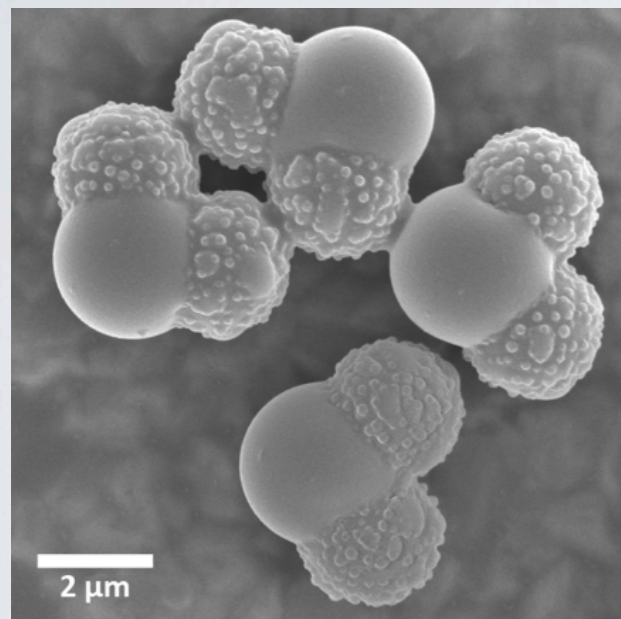
Insufficient steric protection leads to unlimited growth.

Can we design particles that constrain growth in 2D?



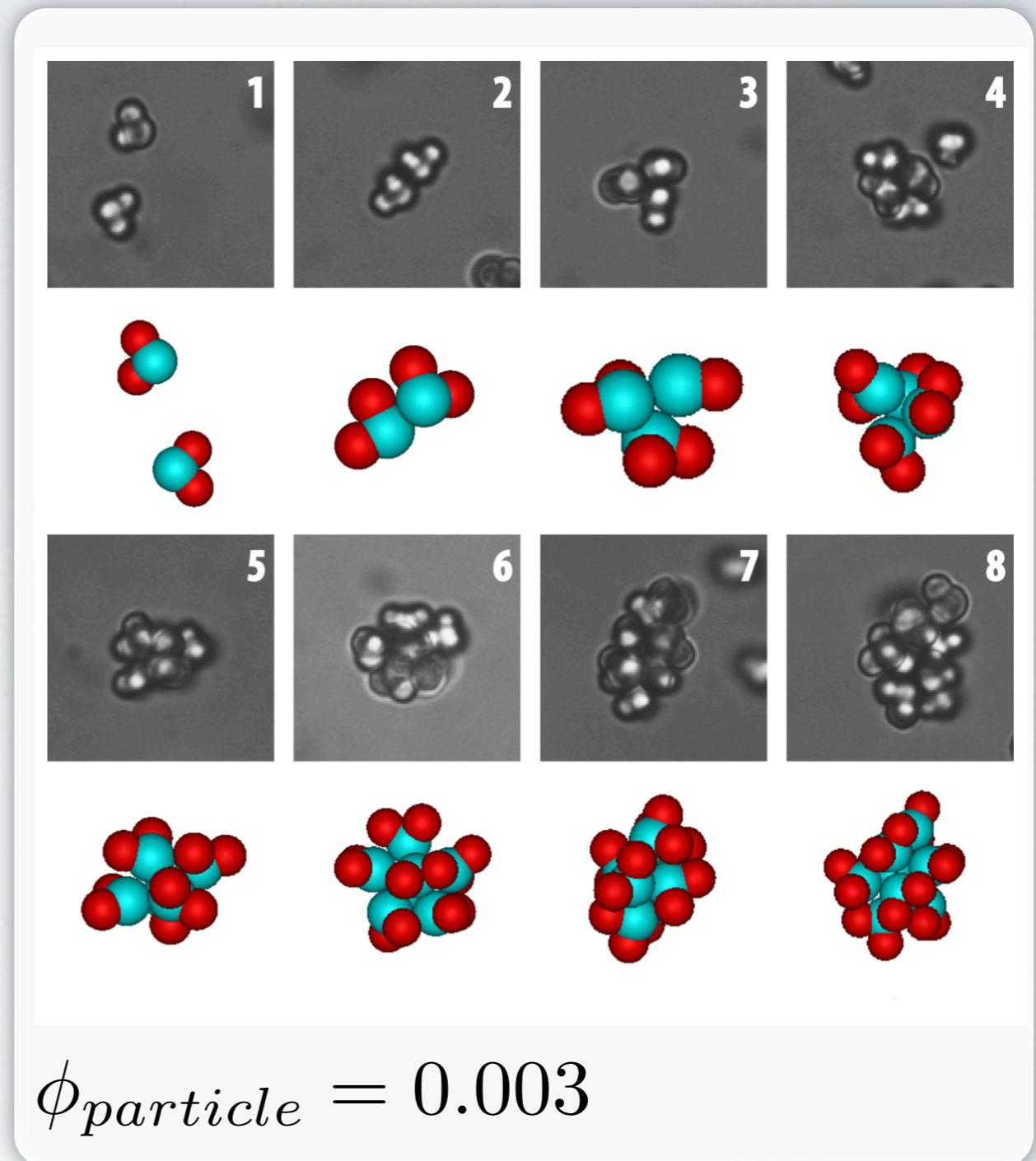
“Mickey Mouse” science

DEPLETION INDUCED ASSEMBLY OF MICKEY MOUSE SHAPED COLLOIDS



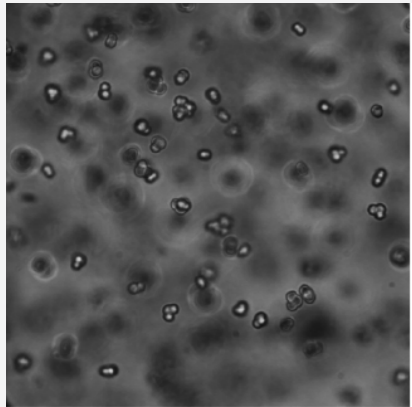
Dextran as depletant ($d=38\text{nm}$)

Experiments headed by J. Wolters,
Simulations by G. Awwisati, T. Vissers
& M. Dijkstra



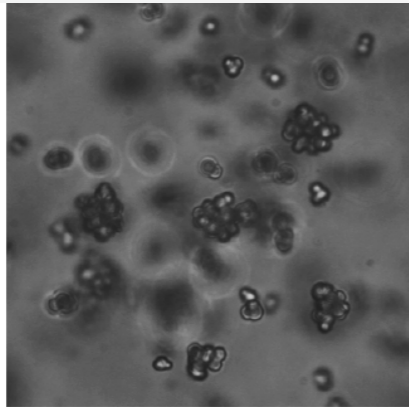
MICKEY MOUSE COLLOIDS ASSEMBLE INTO TUBULAR STRUCTURES

$$\phi_{particle} = 0.01$$



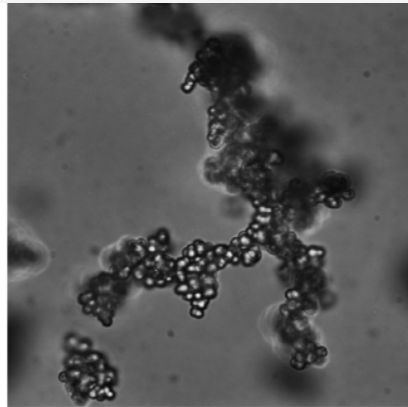
$$\phi_d = 0.20$$

$$\varepsilon = -7k_B T$$



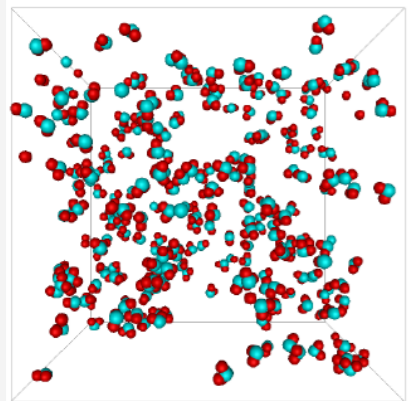
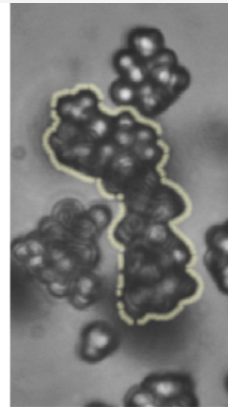
$$\phi_d = 0.21$$

$$\varepsilon = -8k_B T$$

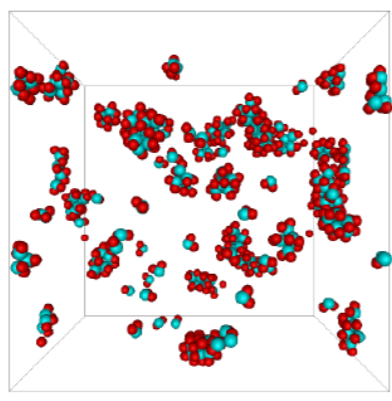


$$\phi_d = 0.24$$

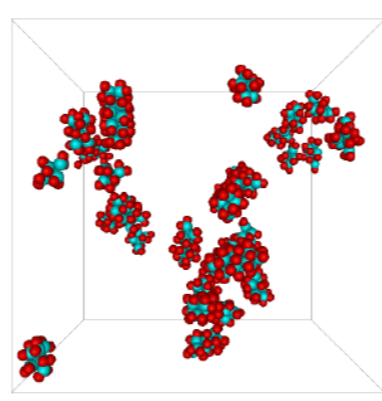
$$\varepsilon = -9k_B T$$



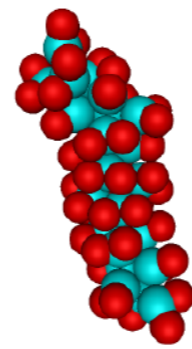
$$\varepsilon = -6k_B T$$



$$\varepsilon = -7k_B T$$

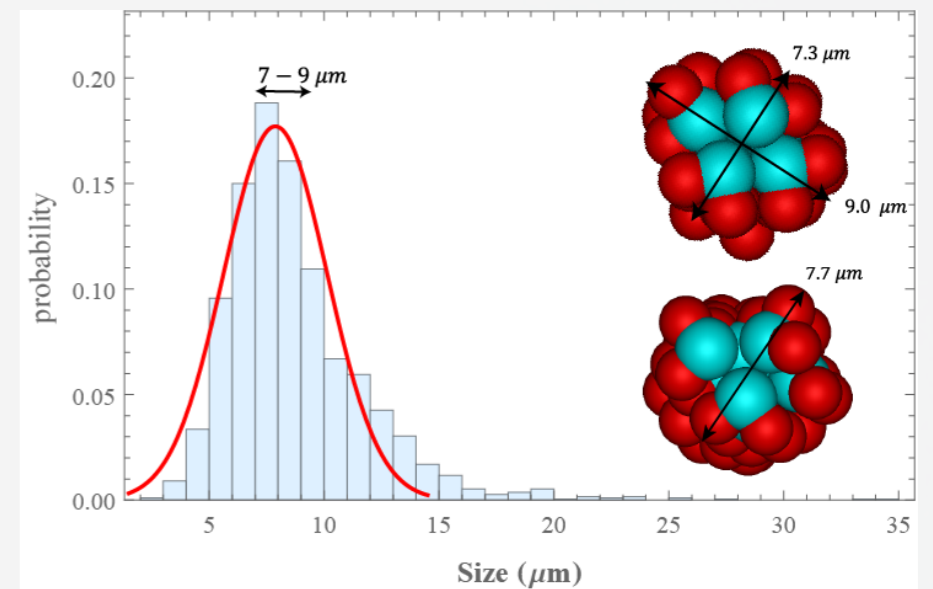


$$\varepsilon = -9k_B T$$



interaction strength \rightarrow

Diameter of tubes



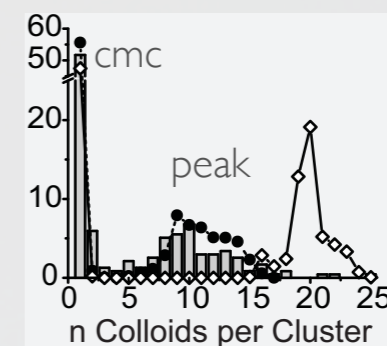
Experiment

Simulations

SUMMARY: SELF-ASSEMBLING ONE-PATCH PARTICLES

Finite size clusters "Colloidal micelle"

Surfactant like behavior!

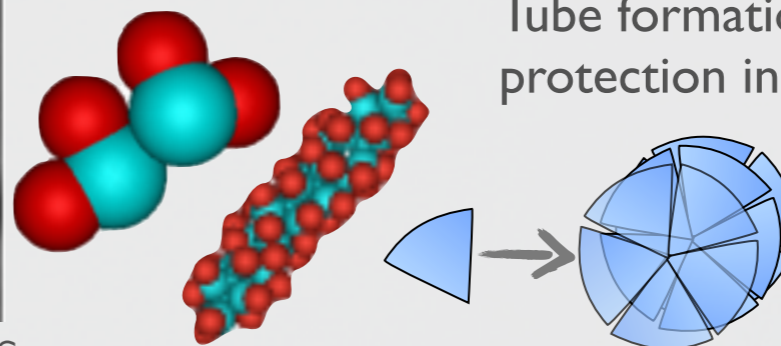


Clusters grow without bounds due to insufficient steric protection.

Kraft et al. PNAS (2012)

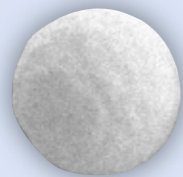
Assembly of tubular structures

Tube formation due to steric protection in 2 dimensions.

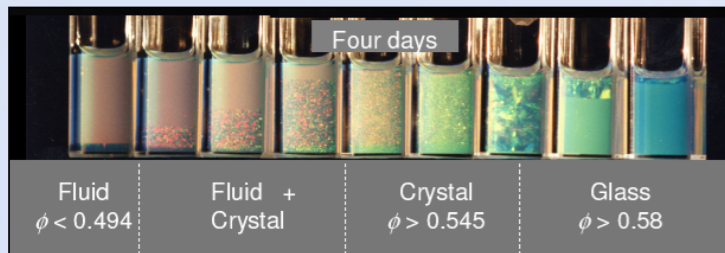


Wolters et al. Soft Matter (2015)

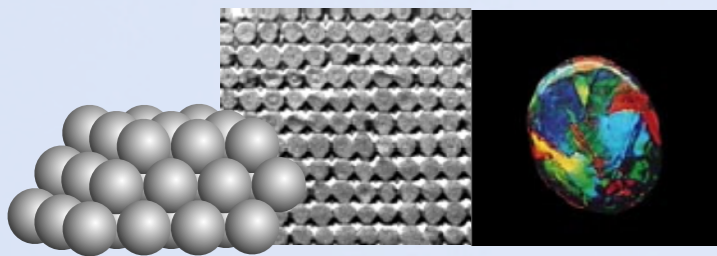
DESIGNED SELF-ASSEMBLY



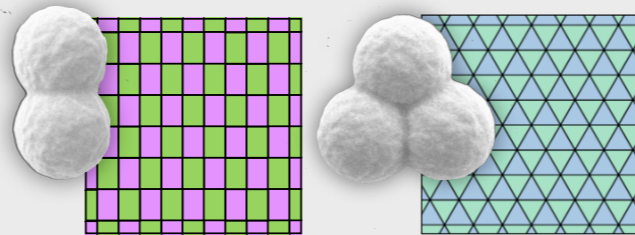
model 'atom'



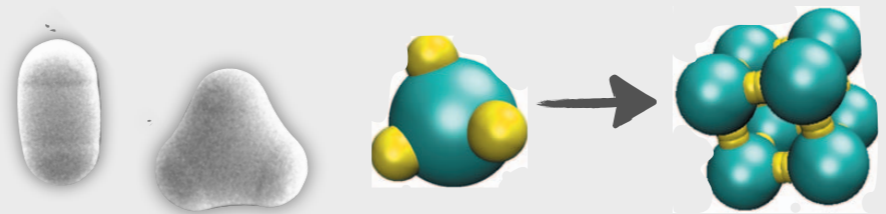
Pusey, van Meegen, Nature (1986)



Anisotropic shape



Anisotropic interactions



Wilber et al. JCP (2009)

Highly specific interactions



+ External guiding rules
+ activity

Particle shape
and interactions



Assembled structure



“Designed self-assembly”

THANK YOU

Leiden University

Vera Meester
Ruben Verweij
Casper van der Wel
Indrani Chakraborty
Hans Frijters
Sabine Matysik
Marcel Winter
Luca Giomi



Experiments

Utrecht University
Willem Kegel
Marlous Kamp
Joost Wolters
Alfons van Blaaderen
Jan Groenewold

Harvard University
Dave Weitz
Kisun Yoon

NYU
David Pine
Stefano Sacanna

Simulations

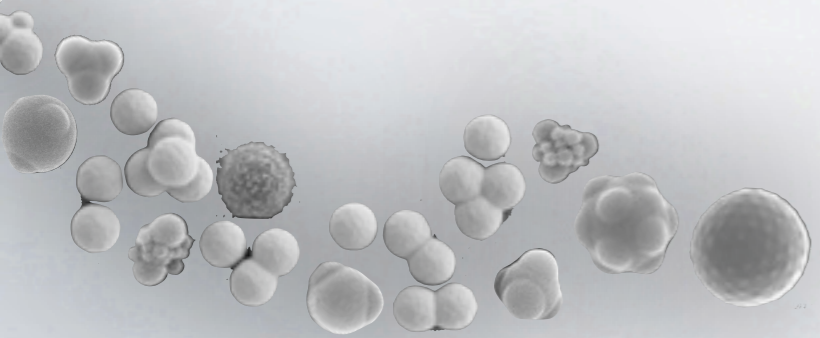
Utrecht University
Ran Ni (now UvA)
Guido Avvisati
Marjolein Dijkstra

University of Edinburgh
Michiel Hermes
Teun Vissers

University of Düsseldorf
Frank Smallenburg

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Rubicon fellowship
VENI grant
Nanofront Gravity grants
DAAD Rise fellowship



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The Netherlands

Thank you for your attention!

